Boundaries, Interfaces, & Duality in 3d SCFTs

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OUTLINE OF THIS TALK

Introduction

Boundary Conditions in 3d $\mathcal{N} = 2$ theories

Dual pairs of boundary conditions

Duality interfaces

(0, 2) theories and the 4-simplex

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MOTIVATION:

- There is more to QFT than just local operators. Extended operators can distinguish QFTs, phases thereof...
- ► Organizing dualities: bulk duality ~ lower dimensional QFTs supported on extended operators
- SUSY + twisting: equalities among local ops/correlation fns in dual theories can admit precise mathematical reformulations leading to interesting (conjectural) equivalences. What about extended operators?

GOALS:

Can we enrich and extend bulk (supersymmetric) dualities by adding (half-BPS) defects?

Today: $3d \mathcal{N} = 2 \text{ IR}$ (Seiberg-like) dualities, codim-1 objects

Given such dual pairs, can we understand the resulting equivalences of underlying mathematical objects?

BOUNDARIES AND INTERFACES

- Interface: domain wall between T, T'
- ► Boundary: "domain wall" between *T*, trivial theory

These objects compose in interesting ways.



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Distinct UV Lagrangian theories that flow to common superconformal fixed point

Simple examples to keep in mind:

• SQED/XYZ: $U(1)_0$, Φ , $\tilde{\Phi} \leftrightarrow W = XYZ$

[Aharony-Hanany-Intriligator-Seiberg-Strassler], [de Boer-Hori-Oz]

► SUSY level-rank: $U(N)_{k+N,k} \leftrightarrow SU(k)_{-k-N}$ [Hsin-Seiberg]

Identify UV boundary conditions $\mathcal{B}, \mathcal{B}^{\vee}$ that flow to the same superconformal \mathcal{B}_{IR}

For simplicity, focus on theories with unitary gauge groups and fundamental/antifundamental matter

Semiclassical bc's in UV YM-CS + matter theories Study (dual pairs of) $3d \mathcal{N} = 2$ theories on a half-space

(expands on [Gadde-Gukov-Putrov])

 $\mathbb{R}^{1,1}\times\mathbb{R}_{\leq 0}$

Half-BPS boundary conditions preserving:

- ► 2d $\mathcal{N} = (0, 2)$ SUSY
- 3d $U(1)_R$ (inherited by 2d superalgebra)

$$\{ Q_+, \bar{Q}_+ \} = -2P_+ \\ \{ Q_-, \bar{Q}_- \} = 2P_- \\ \{ Q_+, \bar{Q}_- \} = -2i(P_\perp - iZ) \\ \{ Q_-, \bar{Q}_+ \} = 2i(P_\perp + iZ)$$



NB: Our UV computations will fail if emergent IR $U(1)_R$!

Building a boundary

Step One:

Write 3d $\mathcal{N} = 2$ multiplets/Lagrangians in 2d $\mathcal{N} = (0,2)$ language \rightarrow

i.e. the 3d theory is a 2d theory on $\mathbb{R}^{1,1}$ whose fields are valued in maps $f:\mathbb{R}\to 3d$ target

•
$$\Phi_{3d} \rightarrow \Phi = \phi + \theta^+ \psi_+ + \dots, \Psi = \bar{\psi}_- + \theta^+ f + \dots$$

• $V_{3d} \rightarrow A = \theta^+ \bar{\theta}^+ A_+, V_- = A_- + \theta^+ \lambda_- + \dots$
 $S = \sigma + iA_\perp - \theta^+ \bar{\lambda}_+ \dots (wz_{gauge})$

Kinetic terms in x_{\perp} repackaged into (0, 2) interactions Fermionic superpotential, "J-term": $\int d\theta^+ \Psi J(\Phi) = \int d\theta^+ \Psi \partial_{\perp} \Phi \rightarrow f \sim \partial_{\perp} \bar{\phi}$

NB: Also E-terms, $\overline{D}_+\Psi = E(\Phi), E.J = 0$ for SUSY

Step Two:

Identify simple reference boundary conditions, then deform:

$$v|_{\partial} = 0 \rightarrow \left(v + \frac{\partial S_{\partial}(u)}{\partial u}\right)|_{\partial} = 0$$

Shortcut: boundary conditions from boundary action e.o.m's:

$$\int_{x^{\perp} \le 0} v \partial_{\perp} u \to \int_{x^{\perp} \le 0} v \partial_{\perp} u + S_{\partial}(u)$$

Reference b.c's: Neumann: $\Psi|_{\partial} = 0 \Rightarrow \partial_{\perp}\phi|_{\partial} = 0, \ \psi_{-}|_{\partial} = 0$

Dirichlet: $\Phi|_{\partial} = 0 \Rightarrow \phi|_{\partial} = 0, \ \psi_+|_{\partial} = 0$

See e.g. [Witten] for general discussion of elliptic boundary conditions Rmk: SUSY^c bdy action $\rightarrow \perp$ comp. of supercurrent total deriv. on ∂ Step One, cont'd: **Reference b.c's for gauge fields: Neumann:** $F_{\perp\pm}|_{\partial} = 0, \sigma|_{\partial} = 0, \lambda_{+}|_{\partial} = 0$ preserves *G* gauge symmetry on boundary (bdy FI term: $\sigma|_{\partial} = t^{\partial}$)

Dirichlet: $A, V_{-}|_{\partial} = 0 \Rightarrow A_{\pm}|_{\partial} = 0, D_{3d} - \partial_{\perp}\sigma|_{\partial} = 0, \lambda_{-}|_{\partial} = 0$ trivializes gauge field at bdy, preserves only *G* global symmetry

Neumann b.c's require projection onto gauge invariants

Dirichlet b.c's admit boundary monopole operators!

Rmk: Can define bdy conditions defined by singular behavior of bulk fields $\rightarrow \partial_i$ simple examples in paper

Step Three:

Couple to extra boundary matter.

Deforms boundary conditions, e.g. Neumann $\rightarrow F_{\perp\pm}|_{\partial} = J_{\pm}^{\partial}$

► Neumann b.c. on chirals in presence of $W_{3d}(\Phi)$: $\int_{x^{\perp} < 0} dx^{\perp} \sum_{a} E_{a} J^{a} = W(\Phi)|_{\partial}$, add 2*d* Fermis Γ s.t.

$$\sum_{i} E^{i}_{\Gamma} J_{i,\Gamma} = -W(\Phi)|_{\partial}$$

choice of matrix factorization

► Neumann b.c. on gauge fields

gauge anomaly cancellation

Dual pairs? Match 't Hooft anomalies and flavor symmetries!

Anomaly Interlude

Claim: A 3d fermion with b.c. $\psi_{\pm}|_{\partial} = 0$ contributes $\pm \frac{1}{2}\mathbf{f}^2$, i.e. *half* the contribution of a purely 2d fermion, to the anomaly polynomial.

• Add a real mass term $im\epsilon^{\alpha\beta}\psi_{\alpha\beta}$

- Flow to IR
- Match UV-IR anomalies

Two effects:

- 1. Integrating out massive fermion generates bckgd CS term at level $\frac{1}{2}$ sign(m)
- 2. Normalizable edge modes (2d fermions) may survive at boundary

Net result:

 $\psi_+|_{\partial} = 0$ has total IR anomaly

$$\frac{1}{2}\operatorname{sign}(m) + \left\{ \begin{array}{cc} 0 & \text{if } m > 0 \\ 1 & \text{if } m < 0 \end{array} \right\} = \frac{1}{2}$$

 $\psi_{-}|_{\partial} = 0$ has total IR anomaly

$$\frac{1}{2}\operatorname{sign}(m) + \left\{ \begin{array}{cc} -1 & \text{if } m > 0 \\ 0 & \text{if } m < 0 \end{array} \right\} = -\frac{1}{2}$$

NB: detailed (but analogous) computation also for gauginos in paper

Rmk: see also [Moore-Freed] for a geometric avatar

Half-Index

$$II_{\mathcal{T},\mathcal{B}}(x;q) = Tr_{Ops_{\mathcal{B}}}(-1)^{R}q^{J+R/2}x^{e}$$

Dual pair of b.c's, same half-index (N gauge b.c: [Gadde-Gukov-Putrov]):

- Path integral on $HS^2 \times S^1$ localization for N gauge b.c.: [Yoshida-Sugiyama]
- ► Operatorial count of Q
 ₊-closed boundary local operators, vacuum character of bdy chiral algebra

$$\begin{split} II_N(x;q) &= \prod_{n \ge 0} \frac{1}{1 - q^n x} =: \frac{1}{(x;q)_\infty} \\ II_D(x;q) &= \prod_{n \ge 0} (1 - q^{n+1} x^{-1}) = (q x^{-1};q)_\infty \end{split}$$

Augment with lines to compute characters of other modules:



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Example: Level-rank duality

Proposal (G simple, k > 0):

- 1. $\mathcal{N} = 2 G_{k+h^{\vee}}$ YMCS with \mathcal{D} b.c. flows to chiral G_k WZW model coupled to $TFT[G_k]$.
- 2. $\mathcal{N} = 2 G_{-k-h^{\vee}}$ YMCS with \mathcal{N} b.c. MUST be coupled to \mathcal{T}_{2d} , flows to \mathcal{T}_{2d}/G_k coupled to $TFT[G_{-k}]$.

 $U(N)_{-k-N,-k} \mathcal{N} + k$ fund. Fermis $\leftrightarrow SU(k)_{k+N} \mathcal{D}$

Half-index of $SU(2)_{1+2}$ with Dirichlet boundary conditions on gauge field

$$II_{\mathcal{T}_1,\mathcal{D}} = \chi_0(SU(2)_1) \tag{1}$$

Half-index of $U(1)_{-2}$ with Neumann bc's + 2 fund. Fermis

$$II_{\mathcal{T}_2,\mathcal{N}} = \chi_0(SU(2)_1) \tag{2}$$

Each equality is nontrivial. Here, equality is a manifestation of level-rank dual *bdy* CFTs! $\frac{2 \text{ fermions}}{U(1)_2} \simeq SU(2)_1$ More generally: $\frac{\text{k fermions}}{U(1)_k} \simeq SU(k)_1$

More Examples

3d \mathcal{N} = 2 dual pairs:

- $\blacktriangleright \ U(N), (N_f, \bar{N}_f) \leftrightarrow U(N_f N), (N_f, \bar{N}_f) + \text{singlets, superpotential [Aharony]}$
- $U(N)_{k+N} \leftrightarrow U(k)_{-k-N}$ (more SUSY level-rank)
- ► free chiral $\leftrightarrow U(1)_{1/2}$ + chiral $\leftrightarrow U(1)_{-1/2}$ + chiral (SUSY particle-vortex)

►
$$U(N)_{k+N-\frac{N_f+N_a}{2}}, (N_f, N_a) \leftrightarrow U(k)_{-k-N+\frac{N_f+N_a}{2}}, (N_a, N_f)$$

[Benini-Closset-Cremonesi]

► $SU(2)_1$ + adjoint \leftrightarrow free chiral [Jafferis-Yin]

In other examples, many dual pairs we find are the simplest possible choices consistent with:

- 1. matching boundary flavor symmetries
- 2. matching boundary 't Hooft anomalies
- 3. gauge anomaly cancellation
- 4. superpotential factorization

BOUNDARY CHIRAL ALGEBRAS

[Costello-Dimofte-Gaiotto]

Dual bdy conditions support \simeq chiral algebras in \bar{Q}_+ -coh

- $\partial_{\perp}, \partial_{\overline{z}}$ exact, but ∂_z is not
- "Holomorphic/topological" twist, see [Aganagic-Costello-McNamara-Vafa]
- $\blacktriangleright Bulk ops \leftrightarrow trivial OPE (\textit{except superpotentials} \rightarrow \textit{quantum corrections})$

SQED: $(NNN) + \Gamma$ dual to XYZ: (NDD) Gauge-invt, \bar{Q}_+ -closed operators with nontrivial OPE? $(\phi\gamma)(z)(\tilde{\phi}\bar{\gamma})(w) \sim \frac{\phi\tilde{\phi}(z)}{(z-w)}$ $(\psi_Y)(z)(\psi_Z)(w) \sim \frac{X(z)}{(z-w)}$

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Build a duality interface



Example

$U(N)_{-k-N,-k}|\Gamma|SU(k)_{k+N}$

Neumann b.c's on either side minimally coupled to Nk free Fermi multiplets Γ (flavor symmetry from Fermis gauged by coupling to RHS)

 \Rightarrow

 Γ form a bifundamental representation of the gauge groups

Can check that total anomaly cancels for k, N > 0

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Geometric interpretation?

[Dimofte-NMP], to appear

 $\begin{array}{l} & \textbf{3d-3d correspondence}_{\text{[Dimofte-Gaiotto-Gukov]}} \\ & \textbf{6d} (2,0) \text{ thy on 3-mfld } M_3 \leftrightarrow \textbf{3d } \mathcal{N} = 2 \text{ theory } \mathcal{T}[M_3] \end{array}$

▶ IR duality of SQED/XYZ \leftrightarrow change of triangulation



 Geometry: Cobordism between two different triangulations = 4-simplex

► Field theory: Change of triangulation implemented by (0,2) duality interface

"Duality" of duality interfaces

Continue this logic:

Different ways of gluing 4-simplex together ↔ Different sequence of duality interfaces



Different gluings realize homeomorphic 4-simplices ... IR equivalence of sequences/collisions of duality

interfaces

Associate bdy chiral algebras to simple 4-manifold operation (+ functional identities from half-indices)

see also [Cecotti-Cordova-Vafa, Nakajima, Vafa-Witten, Gadde-Gukov-Putrov, Dedushenko-Gukov-Putrov,

Nidaiev-Moore, Feigin-Gukov, ...]

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Many interesting adjacent areas!

1. Similar techniques employed for non-SUSY dual pairs

[Gaiotto's Nati-Fest talk, Aitken-Baumgartner-Karch-Robinson, Aitken-Karch-Robinson]

- Consistency of our proposal ← 2d dualities, e.g. (0,2) dualities of Gadde-Gukov-Putrov. Generate new (0, 2) dualities?
- 3. Lift dual b.c's to Seiberg-dual surface operators in 4d $\mathcal{N} = 1$ theories?
- 4. Techniques apply to 3d theories with more SUSY, half-index useful to study mathematical counterparts

[Costello-Gaiotto, Aganagic-Okounkov, ...]

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 String/M-theory construction of 3d N = 4 boundary conditions + applications to geometry, rep theory [WIP w/ Dimofte, Hilburn]

Thank you for your attention!