

Bootstrapping $4d \mathcal{N} = 2$ theories

Madalena Lemos

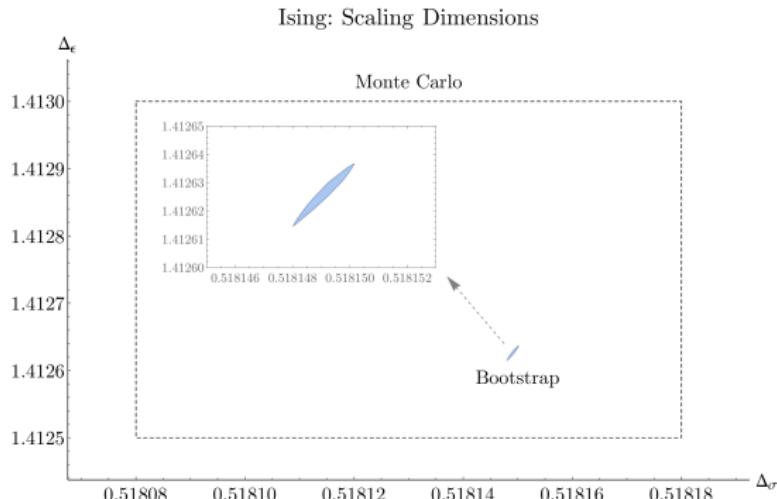


Strings 2018
Jun 28 2018, OIST

together with C. Beem, M. Cornagliotto, P. Liendo, W. Peelaers, L. Rastelli,
V. Schomerus, B. van Rees

3d Ising Model

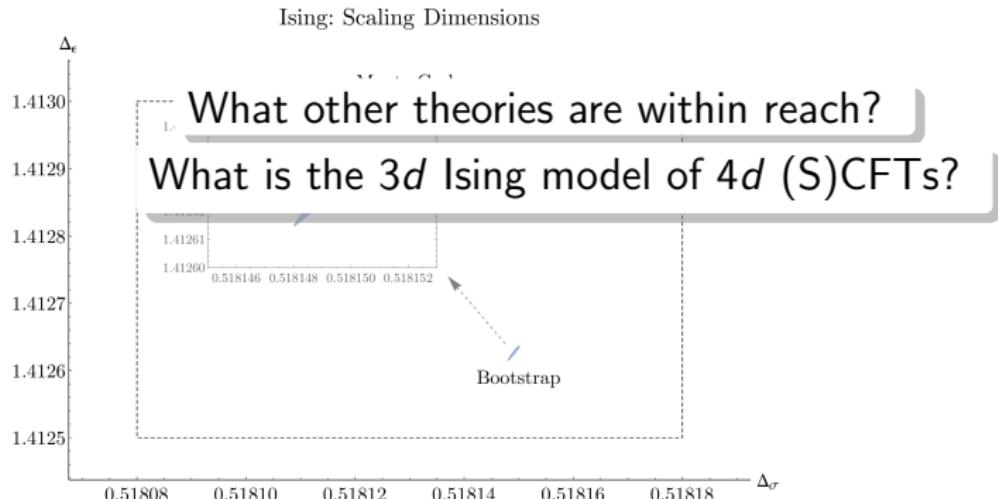
[Poland Simmons-Duffin Kos, Simmons-Duffin, Poland Simmons-Duffin Kos Vichi]



One \mathbb{Z}_2 -even, one \mathbb{Z}_2 -odd relevant scalar operator

3d Ising Model

[Poland Simmons-Duffin Kos, Simmons-Duffin, Poland Simmons-Duffin Kos Vichi]



One \mathbb{Z}_2 -even, one \mathbb{Z}_2 -odd relevant scalar operator

Outline

- ① The Superconformal Bootstrap Program
- ② (A_1, A_2) Argyres-Douglas Theory
- ③ Landscape of $4d \mathcal{N} = 2$ SCFTs
- ④ Summary & Outlook

Outline

- ① The Superconformal Bootstrap Program
- ② (A_1, A_2) Argyres-Douglas Theory
- ③ Landscape of $4d \mathcal{N} = 2$ SCFTs
- ④ Summary & Outlook

The Superconformal Bootstrap Program

What is the space of consistent $4d$ SCFTs?

The Superconformal Bootstrap Program

What is the space of consistent $4d$ SCFTs?

- Maximally supersymmetric theories: $\mathcal{N} = 4$ SYM (?)

The Superconformal Bootstrap Program

What is the space of consistent $4d$ SCFTs?

- Maximally supersymmetric theories: $\mathcal{N} = 4$ SYM (?)
- $\mathcal{N} = 2$ theories: growing list of theories [see Argyres' talk]

The Superconformal Bootstrap Program

What is the space of consistent 4d SCFTs?

- Maximally supersymmetric theories: $\mathcal{N} = 4$ SYM (?)
- $\mathcal{N} = 3$ theories [García-Etxebarria Regalado]
- $\mathcal{N} = 2$ theories: growing list of theories [see Argyres' talk]

The Superconformal Bootstrap Program

What is the space of consistent $4d$ SCFTs?

- Maximally supersymmetric theories: $\mathcal{N} = 4$ SYM (?)
- $\mathcal{N} = 3$ theories [García-Etxebarria Regalado]
- $\mathcal{N} = 2$ theories: growing list of theories [see Argyres' talk]

Can we bootstrap specific theories?

The Superconformal Bootstrap Program

What is the space of consistent 4d SCFTs?

- Maximally supersymmetric theories: $\mathcal{N} = 4$ SYM (?)
- $\mathcal{N} = 3$ theories [García-Etxebarria Regalado]
- $\mathcal{N} = 2$ theories: growing list of theories [see Argyres' talk]

Can we bootstrap specific theories?

- “Simplest” $\mathcal{N} = 2$ Argyres-Douglas theory?

Conformal Bootstrap

Conformal field theory defined by [see Simmons-Duffin's talk]

Set of local operators and *all* their correlation functions

Conformal Bootstrap

Conformal field theory defined by [see Simmons-Duffin's talk]

Set of local operators and *all* their correlation functions

Operator Product Expansion

$$\mathcal{O}_1(x)\mathcal{O}_2(0) = \sum_{k \text{ prim.}} f_{\mathcal{O}_1\mathcal{O}_2\mathcal{O}_k} x^{\Delta_k - \Delta_1 - \Delta_2} (\mathcal{O}_k(0) + \dots)$$

Conformal Bootstrap

Conformal field theory defined by [see Simmons-Duffin's talk]

$$\{\mathcal{O}_{\Delta,\ell,\dots}(x)\} \text{ and } \{f_{\mathcal{O}_i \mathcal{O}_j \mathcal{O}_k}\}$$

Operator Product Expansion

$$\mathcal{O}_1(x) \mathcal{O}_2(0) = \sum_{k \text{ prim.}} f_{\mathcal{O}_1 \mathcal{O}_2 \mathcal{O}_k} x^{\Delta_k - \Delta_1 - \Delta_2} (\mathcal{O}_k(0) + \dots)$$

Conformal Bootstrap

Conformal field theory defined by [see Simmons-Duffin's talk]

$$\{\mathcal{O}_{\Delta,\ell,\dots}(x)\} \text{ and } \{f_{\mathcal{O}_i \mathcal{O}_j \mathcal{O}_k}\}$$

Operator Product Expansion

$$\mathcal{O}_1(x) \mathcal{O}_2(0) = \sum_{k \text{ prim.}} f_{\mathcal{O}_1 \mathcal{O}_2 \mathcal{O}_k} x^{\Delta_k - \Delta_1 - \Delta_2} (\mathcal{O}_k(0) + \dots)$$

Subject to

- ▶ Unitarity
- ▶ Associativity of the operator product algebra

Conformal Bootstrap

Conformal field theory defined by [see Simmons-Duffin's talk]

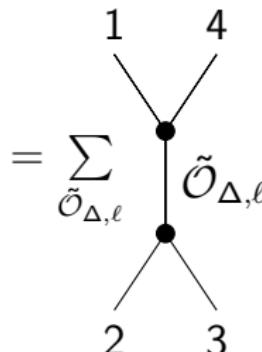
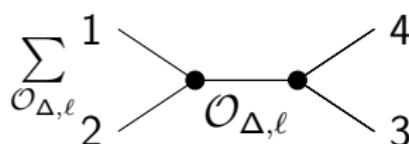
$$\{\mathcal{O}_{\Delta,\ell,\dots}(x)\} \text{ and } \{f_{\mathcal{O}_i \mathcal{O}_j \mathcal{O}_k}\}$$

Operator Product Expansion

$$\mathcal{O}_1(x) \mathcal{O}_2(0) = \sum_{k \text{ prim.}} f_{\mathcal{O}_1 \mathcal{O}_2 \mathcal{O}_k} x^{\Delta_k - \Delta_1 - \Delta_2} (\mathcal{O}_k(0) + \dots)$$

Subject to

- ▶ Unitarity
- ▶ Crossing equations for *all* four-point functions



The Superconformal Bootstrap

Conformal field theory defined by [see Simmons-Duffin's talk]

$$\{\mathcal{O}_{\Delta,\ell,\dots}(x)\} \text{ and } \{f_{\mathcal{O}_i \mathcal{O}_j \mathcal{O}_k}\}$$

The Superconformal Bootstrap

Conformal field theory defined by [see Simmons-Duffin's talk]
 $\{\mathcal{O}_{\Delta,\ell,\dots}(x)\}$ and $\{f_{\mathcal{O}_i \mathcal{O}_j \mathcal{O}_k}\}$

The Superconformal Bootstrap

- ▶ Conformal families \rightsquigarrow Superconformal families

The Superconformal Bootstrap

Conformal field theory defined by [see Simmons-Duffin's talk]
 $\{\mathcal{O}_{\Delta,\ell,\dots}(x)\}$ and $\{f_{\mathcal{O}_i \mathcal{O}_j \mathcal{O}_k}\}$

The Superconformal Bootstrap

- ▶ Conformal families \leadsto Superconformal families
- ▶ Finite re-organization of an infinite amount of data

The Superconformal Bootstrap

Conformal field theory defined by [see Simmons-Duffin's talk]
 $\{\mathcal{O}_{\Delta,\ell,\dots}(x)\}$ and $\{f_{\mathcal{O}_i \mathcal{O}_j \mathcal{O}_k}\}$

The Superconformal Bootstrap

- ▶ Conformal families \rightsquigarrow Superconformal families
- ▶ Finite re-organization of an infinite amount of data

Q: Is there a solvable truncation of the crossing equations?

The Superconformal Bootstrap

Conformal field theory defined by [see Simmons-Duffin's talk]
 $\{\mathcal{O}_{\Delta,\ell,\dots}(x)\}$ and $\{f_{\mathcal{O}_i \mathcal{O}_j \mathcal{O}_k}\}$

The Superconformal Bootstrap

- ▶ Conformal families \rightsquigarrow Superconformal families
- ▶ Finite re-organization of an infinite amount of data

Q: Is there a solvable truncation of the crossing equations?

- Yes, for $4d \mathcal{N} \geq 2$ [Beem ML Liendo Peelaers Rastelli van Rees]
(and also $6d \mathcal{N} = (2,0)$ and $2d \mathcal{N} = (0,4)$ [Beem Rastelli van Rees])

The Superconformal Bootstrap

Conformal field theory defined by [see Simmons-Duffin's talk]
 $\{\mathcal{O}_{\Delta,\ell,\dots}(x)\}$ and $\{f_{\mathcal{O}_i \mathcal{O}_j \mathcal{O}_k}\}$

The Superconformal Bootstrap

- ▶ Conformal families \rightsquigarrow Superconformal families
- ▶ Finite re-organization of an infinite amount of data

Q: Is there a solvable truncation of the crossing equations?

- Yes, for $4d \mathcal{N} \geq 2$ [Beem ML Liendo Peelaers Rastelli van Rees]
(and also $6d \mathcal{N} = (2,0)$ and $2d \mathcal{N} = (0,4)$ [Beem Rastelli van Rees])
- Subsector $\mathcal{N} \geq 2$ SCFTs captured by $2d$ chiral algebra

A solvable subsector

4d $\mathcal{N} = 2$ SCFTs \rightarrow 2d chiral algebra

A solvable subsector

4d $\mathcal{N} = 2$ SCFTs \rightarrow 2d chiral algebra

- ▶ $SU(2)_R$ current \mapsto 2d stress tensor $T(z)$

A solvable subsector

4d $\mathcal{N} = 2$ SCFTs \rightarrow 2d chiral algebra

- ▶ $\underbrace{SU(2)_R}_{\in \text{Super-stress tensor multiplet}}$ current \mapsto 2d stress tensor $T(z)$

A solvable subsector

$4d \mathcal{N} \geq 2$ SCFTs \rightarrow $2d$ chiral algebra

- ▶ Super-stress tensor multiplet $_{4d}$ \mapsto (Super-)stress tensor $_{2d}$

A solvable subsector

$4d \mathcal{N} \geq 2$ SCFTs \rightarrow $2d$ chiral algebra

- ▶ Super-stress tensor multiplet $_{4d}$ \mapsto (Super-)stress tensor $_{2d}$

A trivial statement in $2d$

- (super-)stress tensor four-point function fixed in terms of c_{2d}

A solvable subsector

4d $\mathcal{N} \geq 2$ SCFTs \rightarrow 2d chiral algebra

- ▶ Super-stress tensor multiplet_{4d} \mapsto (Super-)stress tensor_{2d}

A trivial statement in 2d

- (super-)stress tensor four-point function fixed in terms of
 c_{2d} ($\langle TT \rangle \propto c$)

A solvable subsector

$4d \mathcal{N} \geq 2$ SCFTs \rightarrow $2d$ chiral algebra

- ▶ Super-stress tensor multiplet $_{4d}$ \mapsto (Super-)stress tensor $_{2d}$

A trivial statement in $2d$

- (super-)stress tensor four-point function fixed in terms of
 $c_{2d} = -12c_{4d}$ ($\langle TT \rangle \propto c$)

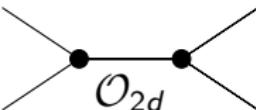
A solvable subsector

4d $\mathcal{N} \geq 2$ SCFTs \rightarrow 2d chiral algebra

- ▶ Super-stress tensor multiplet_{4d} \mapsto (Super-)stress tensor_{2d}

A trivial statement in 2d

- (super-)stress tensor four-point function fixed in terms of
 $c_{2d} = -12c_{4d}$ ($\langle TT \rangle \propto c$)
- 2d Superblock decomposition:

$$\sum_{\mathcal{O}_{2d}} f_{\mathcal{O}_{2d}}^2$$


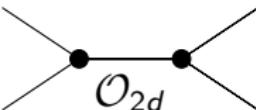
A solvable subsector

4d $\mathcal{N} \geq 2$ SCFTs \rightarrow 2d chiral algebra

- ▶ Super-stress tensor multiplet_{4d} \mapsto (Super-)stress tensor_{2d}

A trivial statement in 2d

- (super-)stress tensor four-point function fixed in terms of
 $c_{2d} = -12c_{4d}$ ($\langle TT \rangle \propto c$)
- 2d Superblock decomposition:

$$\sum_{\mathcal{O}_{2d}} f_{\mathcal{O}_{2d}}^2$$


- $f_{\mathcal{O}_{2d}}^2$

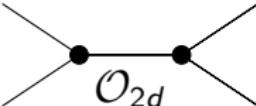
A solvable subsector

4d $\mathcal{N} \geq 2$ SCFTs \rightarrow 2d chiral algebra

- ▶ Super-stress tensor multiplet_{4d} \mapsto (Super-)stress tensor_{2d}

A trivial statement in 2d

- (super-)stress tensor four-point function fixed in terms of
 $c_{2d} = -12c_{4d}$ ($\langle TT \rangle \propto c$)
- 2d Superblock decomposition:

$$\sum_{\mathcal{O}_{2d}} f_{\mathcal{O}_{2d}}^2$$


$$\rightarrow f_{\mathcal{O}_{2d}}^2 \rightsquigarrow f_{\mathcal{O}_{4d}}^2$$

assumptions: interacting theory, unique stress tensor

A solvable subsector

4d $\mathcal{N} \geq 2$ SCFTs \rightarrow 2d chiral algebra

- ▶ Super-stress tensor multiplet_{4d} \mapsto (Super-)stress tensor_{2d}

A trivial statement in 2d

- (super-)stress tensor four-point function fixed in terms of
 $c_{2d} = -12c_{4d}$ ($\langle TT \rangle \propto c$)
- 2d Superblock decomposition:

$$\sum_{\mathcal{O}_{2d}} f_{\mathcal{O}_{2d}}^2$$

$$f_{\mathcal{O}_{2d}}^2 \rightsquigarrow f_{\mathcal{O}_{4d}}^2 \underbrace{\geqslant}_{4d \text{ unitarity}} 0$$

assumptions: interacting theory, unique stress tensor

A solvable subsector

4d $\mathcal{N} \geq 2$ SCFTs \rightarrow 2d chiral algebra

- ▶ Super-stress tensor multiplet_{4d} \mapsto (Super-)stress tensor_{2d}

A trivial statement in 2d

- (super-)stress tensor four-point function fixed in terms of
 $c_{2d} = -12c_{4d}$ ($\langle TT \rangle \propto c$)
- 2d Superblock decomposition:

$$\sum_{\mathcal{O}_{2d}} f_{\mathcal{O}_{2d}}^2$$

$$\rightarrow f_{\mathcal{O}_{2d}}^2 \rightsquigarrow f_{\mathcal{O}_{4d}}^2 \underbrace{\geqslant}_{4d \text{ unitarity}} 0 \Rightarrow \text{New unitarity bounds}$$

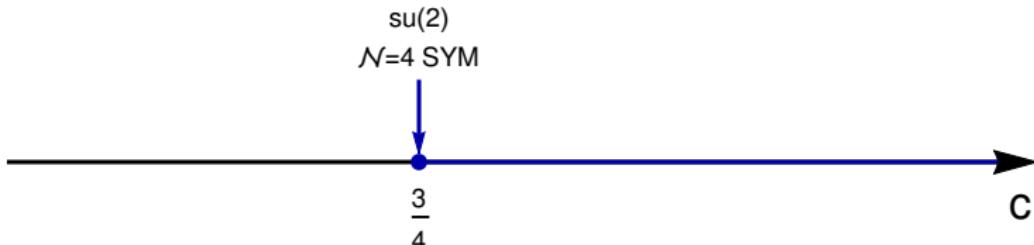
assumptions: interacting theory, unique stress tensor

Landscape of $4d \mathcal{N} \geq 2$ SCFTs

From $2d$ (super-)stress tensor four-point function

(assumptions: interacting theory, unique stress tensor)

→ $4d \mathcal{N} = 4$ SCFTs $c = a \geq \frac{3}{4}$ [Beem Rastelli van Rees]



Landscape of 4d $\mathcal{N} \geq 2$ SCFTs

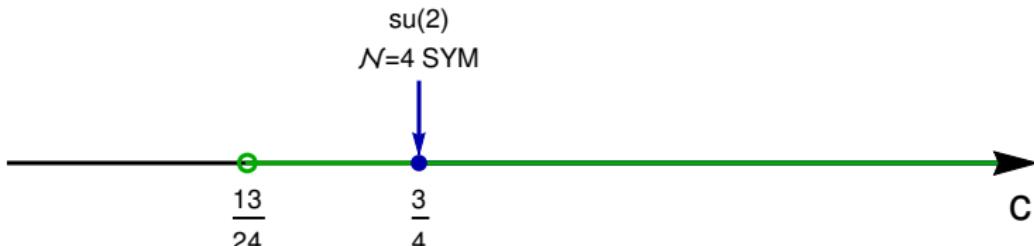
From 2d (super-)stress tensor four-point function

(assumptions: interacting theory, unique stress tensor)

→ 4d $\mathcal{N} = 4$ SCFTs $c = a \geq \frac{3}{4}$ [Beem Rastelli van Rees]

→ 4d $\mathcal{N} \geq 3$ SCFTs $c = a > \frac{13}{24}$ [Cornagliotto ML Schomerus]

from interpreting \mathcal{O}_{2d} as a 4d operator



Landscape of $4d \mathcal{N} \geq 2$ SCFTs

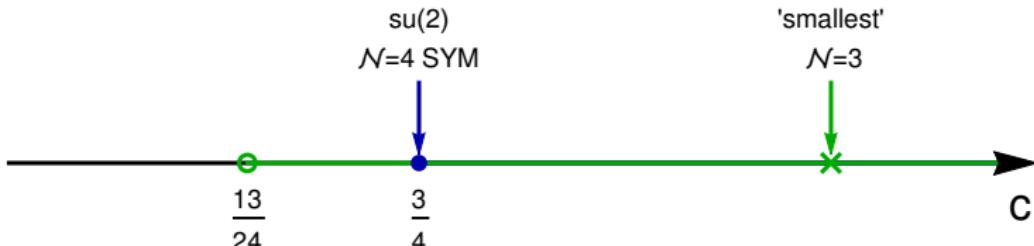
From $2d$ (super-)stress tensor four-point function

(assumptions: interacting theory, unique stress tensor)

$\rightarrow 4d \mathcal{N} = 4$ SCFTs $c = a \geq \frac{3}{4}$ [Beem Rastelli van Rees]

$\rightarrow 4d \mathcal{N} \geq 3$ SCFTs $c = a > \frac{13}{24}$ [Cornagliotto ML Schomerus]

from interpreting \mathcal{O}_{2d} as a $4d$ operator



Landscape of 4d $\mathcal{N} \geq 2$ SCFTs

From 2d (super-)stress tensor four-point function

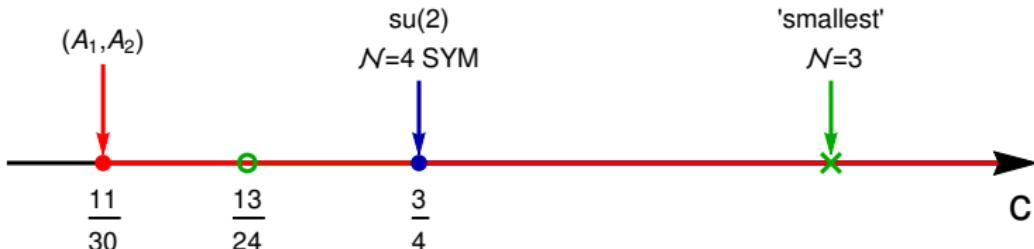
(assumptions: interacting theory, unique stress tensor)

→ 4d $\mathcal{N} = 4$ SCFTs $c = a \geq \frac{3}{4}$ [Beem Rastelli van Rees]

→ 4d $\mathcal{N} \geq 3$ SCFTs $c = a > \frac{13}{24}$ [Cornagliotto ML Schomerus]

from interpreting \mathcal{O}_{2d} as a 4d operator

→ 4d $\mathcal{N} \geq 2$ SCFTs $c \geq \frac{11}{30}$ [Liendo Ramirez Seo]



Landscape of 4d $\mathcal{N} \geq 2$ SCFTs

From 2d (super-)stress tensor four-point function

(assumptions: interacting theory, unique stress tensor)

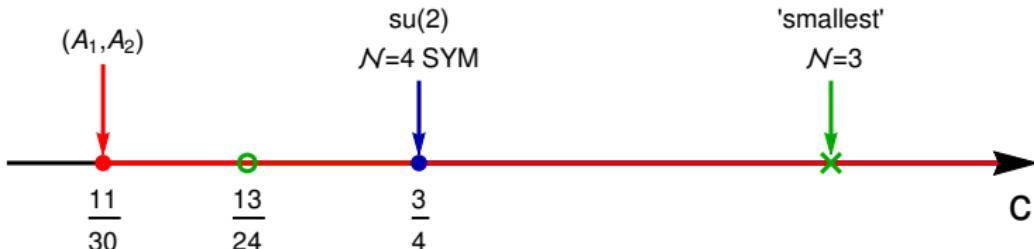
→ 4d $\mathcal{N} = 4$ SCFTs $c = a \geq \frac{3}{4}$ [Beem Rastelli van Rees]

→ 4d $\mathcal{N} \geq 3$ SCFTs $c = a > \frac{13}{24}$ [Cornagliotto ML Schomerus]

from interpreting \mathcal{O}_{2d} as a 4d operator

→ 4d $\mathcal{N} \geq 2$ SCFTs $c \geq \frac{11}{30}$ [Liendo Ramirez Seo]

↪ Saturated by the (A_1, A_2) Argyres-Douglas theory



Outline

- ① The Superconformal Bootstrap Program
- ② (A_1, A_2) Argyres-Douglas Theory
- ③ Landscape of $4d \mathcal{N} = 2$ SCFTs
- ④ Summary & Outlook

The “simplest” Argyres-Douglas theory

- Originally obtained on the Coulomb branch of a 4d $\mathcal{N} = 2$ susy gauge theory with gauge group $SU(3)$

The “simplest” Argyres-Douglas theory

- Originally obtained on the Coulomb branch of a 4d $\mathcal{N} = 2$ susy gauge theory with gauge group $SU(3)$
- $\mathcal{N} = 1$ Lagrangian description [see Song's talk]

The “simplest” Argyres-Douglas theory

- Originally obtained on the Coulomb branch of a 4d $\mathcal{N} = 2$ susy gauge theory with gauge group $SU(3)$
- $\mathcal{N} = 1$ Lagrangian description [see Song's talk]
- Strongly coupled isolated SCFT – no marginal deformations

The “simplest” Argyres-Douglas theory

- Originally obtained on the Coulomb branch of a 4d $\mathcal{N} = 2$ susy gauge theory with gauge group $SU(3)$
- $\mathcal{N} = 1$ Lagrangian description [see Song's talk]
- Strongly coupled isolated SCFT – no marginal deformations
- Just another SCFT

The “simplest” Argyres-Douglas theory

- Originally obtained on the Coulomb branch of a 4d $\mathcal{N} = 2$ susy gauge theory with gauge group $SU(3)$
- $\mathcal{N} = 1$ Lagrangian description [see Song's talk]
- Strongly coupled isolated SCFT – no marginal deformations
- Just another SCFT
- Chiral algebra $[(A_1, A_2)]$ = Lee-Yang minimal model
[Beem Rastelli]

The “simplest” Argyres-Douglas theory

- Originally obtained on the Coulomb branch of a 4d $\mathcal{N} = 2$ susy gauge theory with gauge group $SU(3)$
- $\mathcal{N} = 1$ Lagrangian description [see Song's talk]
- Strongly coupled isolated SCFT – no marginal deformations
- Just another SCFT
- Chiral algebra $[(A_1, A_2)]$ = Lee-Yang minimal model
[Beem Rastelli]

Our tools beyond protected subsector

- ▶ Numerical bootstrap
 - [Rattazzi Rychkov Tonni Vichi]

The “simplest” Argyres-Douglas theory

- Originally obtained on the Coulomb branch of a 4d $\mathcal{N} = 2$ susy gauge theory with gauge group $SU(3)$
- $\mathcal{N} = 1$ Lagrangian description [see Song's talk]
- Strongly coupled isolated SCFT – no marginal deformations
- Just another SCFT
- Chiral algebra $[(A_1, A_2)]$ = Lee-Yang minimal model
[Beem Rastelli]

Our tools beyond protected subsector

- ▶ Numerical bootstrap
 - [Rattazzi Rychkov Tonni Vichi]
- ▶ Lightcone bootstrap
 - [Fitzpatrick Kaplan Poland Simmons-Duffin, Komargodski Zhiboedov]

The “simplest” Argyres-Douglas theory

- Originally obtained on the Coulomb branch of a 4d $\mathcal{N} = 2$ susy gauge theory with gauge group $SU(3)$
- $\mathcal{N} = 1$ Lagrangian description [see Song's talk]
- Strongly coupled isolated SCFT – no marginal deformations
- Just another SCFT
- Chiral algebra $[(A_1, A_2)] = \text{Lee-Yang minimal model}$
[Beem Rastelli]

Our tools beyond protected subsector

- ▶ Numerical bootstrap
 - [Rattazzi Rychkov Tonni Vichi]
- ▶ Lightcone bootstrap
 - [Fitzpatrick Kaplan Poland Simmons-Duffin, Komargodski Zhiboedov]
 - Lorentzian inversion formula of [Caron-Huot]

The “simplest” Argyres-Douglas theory

How can we approach it?

The “simplest” Argyres-Douglas theory

How can we approach it?

- ▶ Known: 4d $\mathcal{N} = 2$ chiral operator ϕ

$$\Delta_\phi = \frac{6}{5}$$

The “simplest” Argyres-Douglas theory

How can we approach it?

- ▶ Known: 4d $\mathcal{N} = 2$ chiral operator ϕ $(\mathcal{Q}_\alpha^I \phi = 0)$

$$\Delta_\phi = \frac{6}{5}$$

The “simplest” Argyres-Douglas theory

How can we approach it?

- ▶ Known: 4d $\mathcal{N} = 2$ chiral operator ϕ $(\mathcal{Q}_\alpha^I \phi = 0)$

$$\Delta_\phi = \frac{6}{5}$$

$U(1)_r$ charge $r = \Delta_\phi$

The “simplest” Argyres-Douglas theory

How can we approach it?

- ▶ Known: 4d $\mathcal{N} = 2$ chiral operator ϕ $(\mathcal{Q}_\alpha^I \phi = 0)$
$$\Delta_\phi = \frac{6}{5}$$
 $U(1)_r$ charge $r = \Delta_\phi$
- ▶ Study $\langle \phi(x_1) \phi(x_2) \bar{\phi}(x_3) \bar{\phi}(x_4) \rangle$

The “simplest” Argyres-Douglas theory

How can we approach it?

- Known: 4d $\mathcal{N} = 2$ chiral operator ϕ ($\mathcal{Q}_\alpha^I \phi = 0$)

$$\Delta_\phi = \frac{6}{5}$$

$U(1)_r$ charge $r = \Delta_\phi$

- Study $\langle \phi(x_1) \phi(x_2) \underbrace{\bar{\phi}(x_3)}_{\text{conjugate of } \phi} \bar{\phi}(x_4) \rangle$

The “simplest” Argyres-Douglas theory

How can we approach it?

- ▶ Known: $4d \mathcal{N} = 2$ chiral operator ϕ ($\mathcal{Q}_\alpha^I \phi = 0$)
$$\boxed{\Delta_\phi = \frac{6}{5}}$$
$$U(1)_r \text{ charge } r = \Delta_\phi$$

- ▶ Study $\langle \phi(x_1) \phi(x_2) \underbrace{\bar{\phi}(x_3)}_{\text{conjugate of } \phi} \bar{\phi}(x_4) \rangle$

- ▶ Two OPE channels:

$$\hookrightarrow \phi\phi \sim \phi^2 + \dots$$

The “simplest” Argyres-Douglas theory

How can we approach it?

- ▶ Known: 4d $\mathcal{N} = 2$ chiral operator ϕ ($\mathcal{Q}_\alpha^I \phi = 0$)
$$\boxed{\Delta_\phi = \frac{6}{5}}$$
$$U(1)_r \text{ charge } r = \Delta_\phi$$
- ▶ Study $\langle \phi(x_1) \phi(x_2) \underbrace{\bar{\phi}(x_3)}_{\text{conjugate of } \phi} \bar{\phi}(x_4) \rangle$
- ▶ Two OPE channels:
 - ↪ $\phi\phi \sim \phi^2 + \dots$
 - ↪ $\phi\bar{\phi} \sim \text{Identity} + \text{Super-stress tensor} + \dots$

The “simplest” Argyres-Douglas theory

How can we approach it?

- Known: 4d $\mathcal{N} = 2$ chiral operator ϕ ($\mathcal{Q}_\alpha^I \phi = 0$)

$$\Delta_\phi = \frac{6}{5}$$

$U(1)_r$ charge $r = \Delta_\phi$

- Study $\langle \phi(x_1) \phi(x_2) \underbrace{\bar{\phi}(x_3)}_{\text{conjugate of } \phi} \bar{\phi}(x_4) \rangle$

- Two OPE channels:

$$\hookrightarrow \phi\phi \sim \phi^2 + \dots$$

$$\hookrightarrow \phi\bar{\phi} \sim \text{Identity} + \text{Super-stress tensor} + \dots$$

- Conformal blocks \leadsto superconformal blocks

The “simplest” Argyres-Douglas theory

How can we approach it?

- ▶ Known: $4d \mathcal{N} = 2$ chiral operator ϕ ($\mathcal{Q}_\alpha^I \phi = 0$)
$$\Delta_\phi = \frac{6}{5}$$
$$U(1)_r \text{ charge } r = \Delta_\phi$$

- ▶ Study $\langle \phi(x_1) \phi(x_2) \underbrace{\bar{\phi}(x_3)}_{\text{conjugate of } \phi} \bar{\phi}(x_4) \rangle$

- ▶ Two OPE channels:

$$\hookrightarrow \phi\phi \sim \phi^2 + \dots$$

$$\hookrightarrow \phi\bar{\phi} \sim \text{Identity} + \text{Super-stress tensor} + \dots$$

- ▶ Conformal blocks \rightsquigarrow superconformal blocks

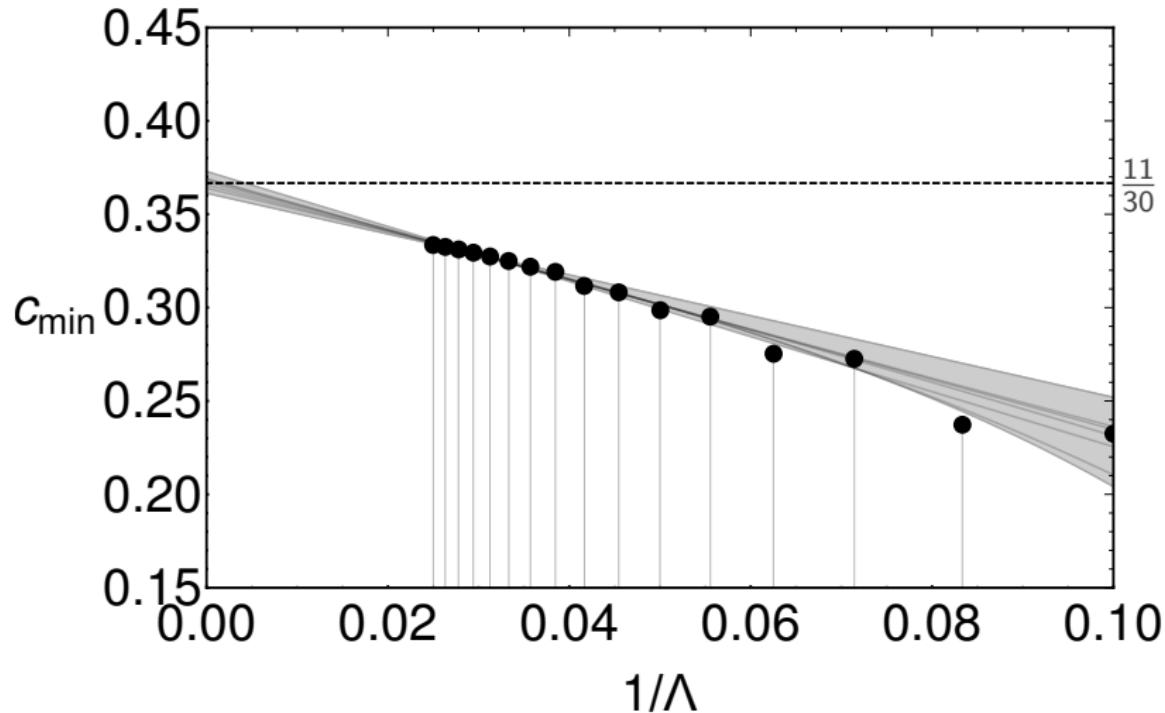
(only in $\phi\bar{\phi}$ channel) [Fitzpatrick Kaplan Khandker Li Poland Simmons-Duffin]

Minimum allowed central charge

Does $\langle \phi\bar{\phi}\bar{\phi}\bar{\phi} \rangle$ know about $c \geq \frac{11}{30}$?

Minimum allowed central charge

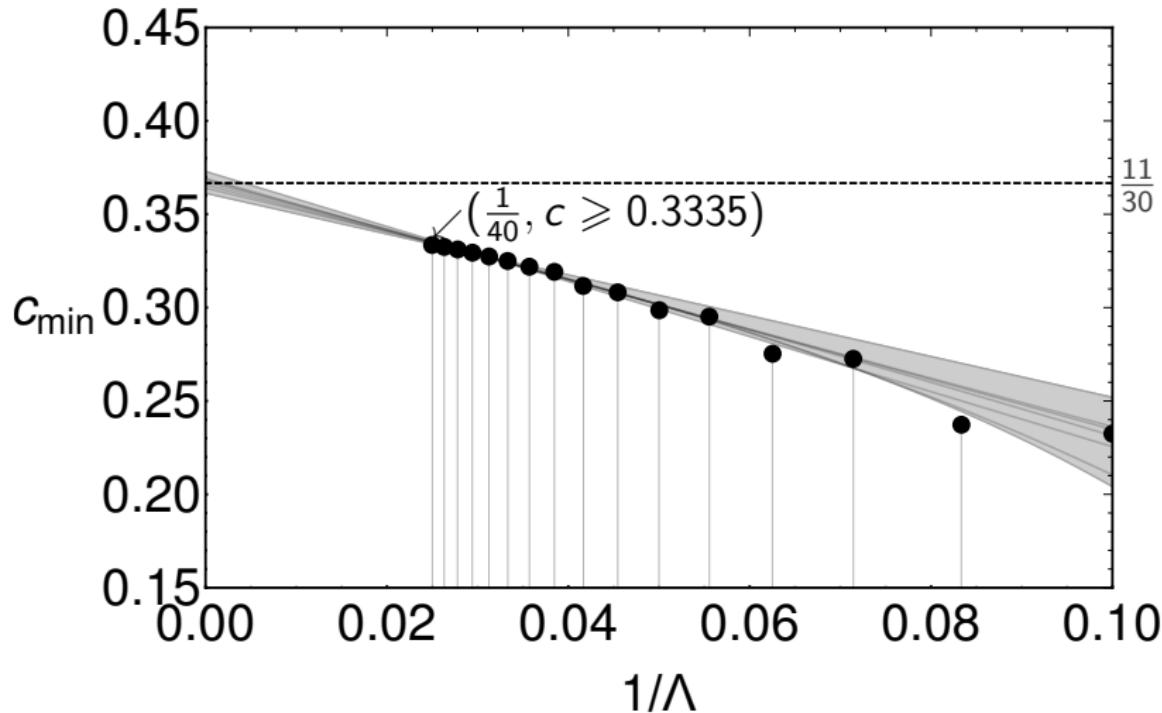
Does $\langle \phi \bar{\phi} \bar{\phi} \bar{\phi} \rangle$ know about $c \geq \frac{11}{30}$?



[Cornagliotto ML Liendo]

Minimum allowed central charge

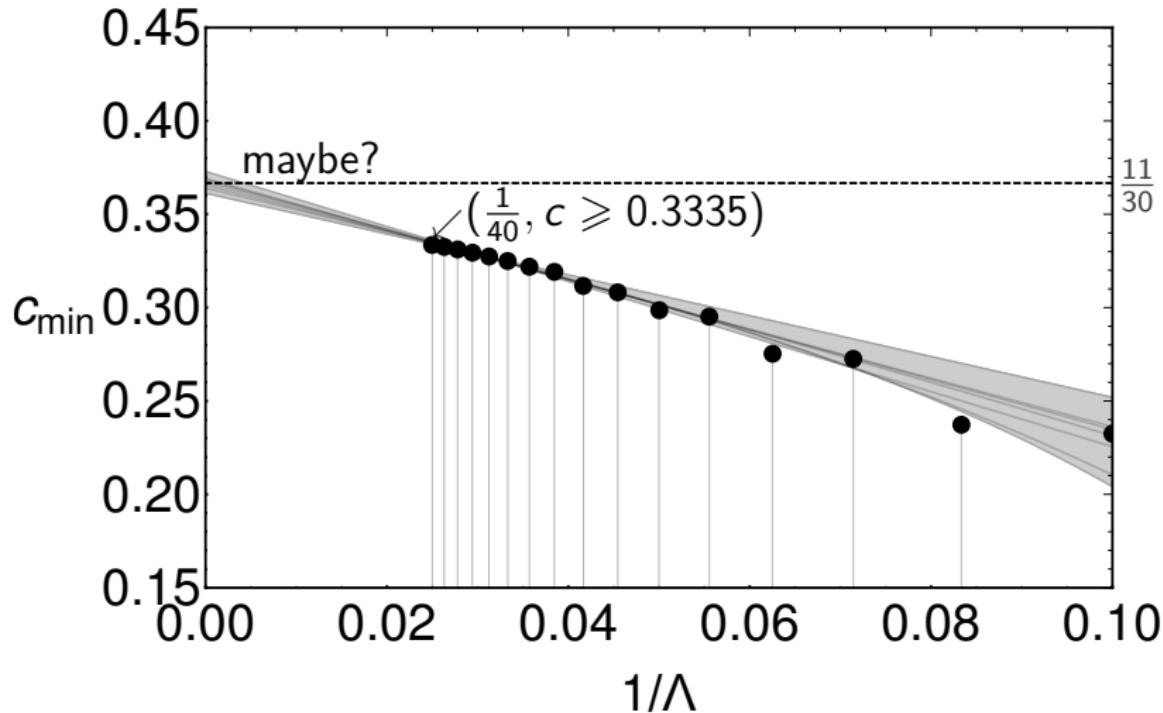
Does $\langle \phi \bar{\phi} \bar{\phi} \bar{\phi} \rangle$ know about $c \geq \frac{11}{30}$?



[Cornagliotto ML Liendo]

Minimum allowed central charge

Does $\langle \phi \bar{\phi} \bar{\phi} \bar{\phi} \rangle$ know about $c \geq \frac{11}{30}$?



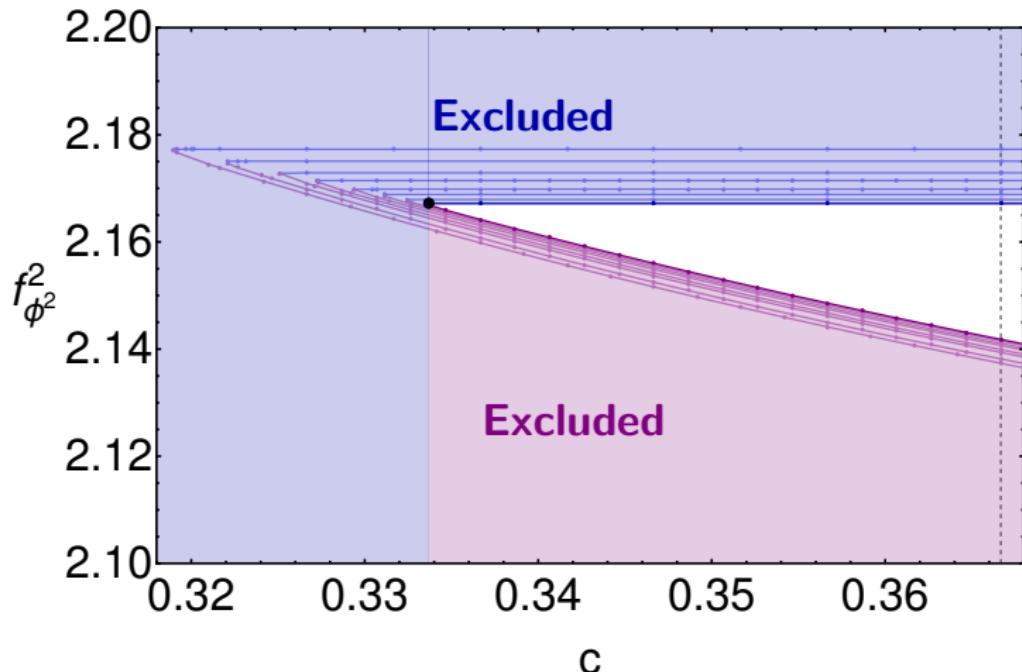
[Cornagliotto ML Liendo]

Bounding OPE coefficients

$$\phi\phi \sim \underbrace{f_{\phi^2}^2}_{\text{unknown}} \underbrace{\phi^2}_{\Delta=2\Delta_\phi} + \dots$$

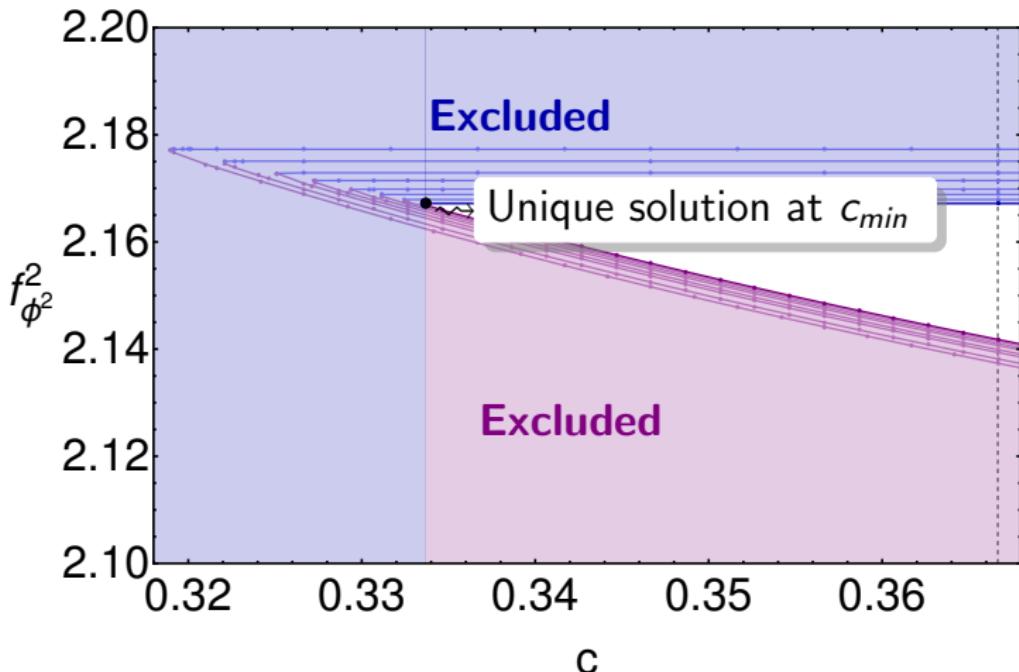
Bounding OPE coefficients

$$\phi\phi \sim \underbrace{f_{\phi^2}^2}_{\text{unknown}} \underbrace{\phi^2}_{\Delta=2\Delta_\phi} + \dots$$



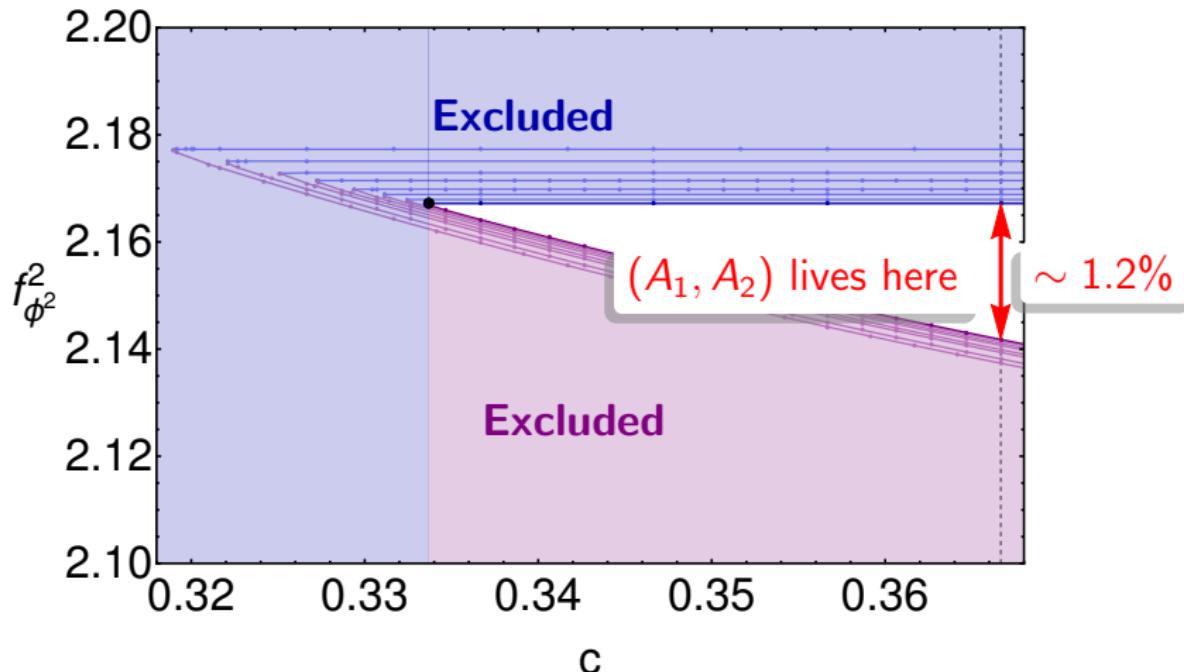
Bounding OPE coefficients

$$\phi\phi \sim \underbrace{f_\phi^2}_{\text{unknown}} \underbrace{\phi^2}_{\Delta=2\Delta_\phi} + \dots$$



Bounding OPE coefficients

$$\phi\phi \sim \underbrace{f_\phi^2}_{\text{unknown}} \underbrace{\phi^2}_{\Delta=2\Delta_\phi} + \dots$$

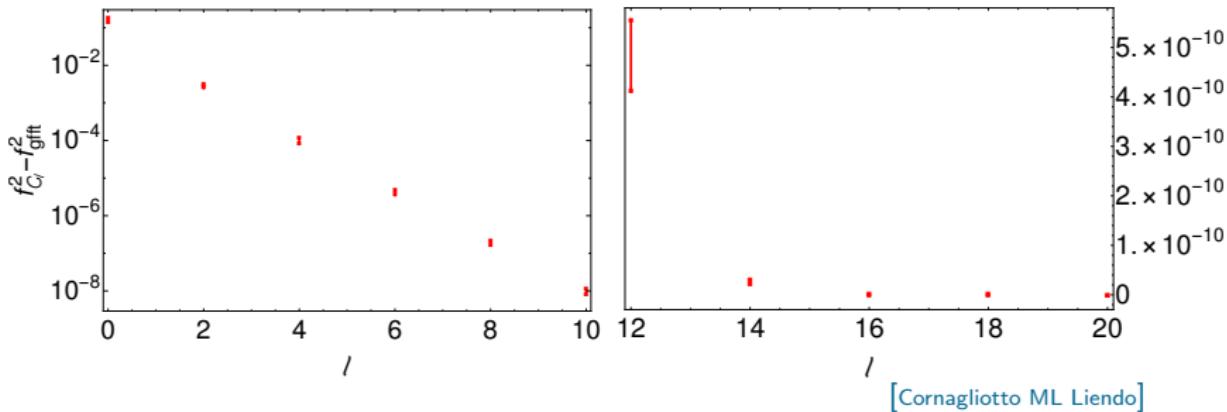


Lorentzian inversion formula

$$\phi\phi \sim f_{\phi^2}^2 \underbrace{\phi^2}_{\Delta=2\Delta_\phi} + f_{C_\ell}^2 \underbrace{C_{\ell>0}}_{\Delta=2\Delta_\phi+\ell} + \dots$$

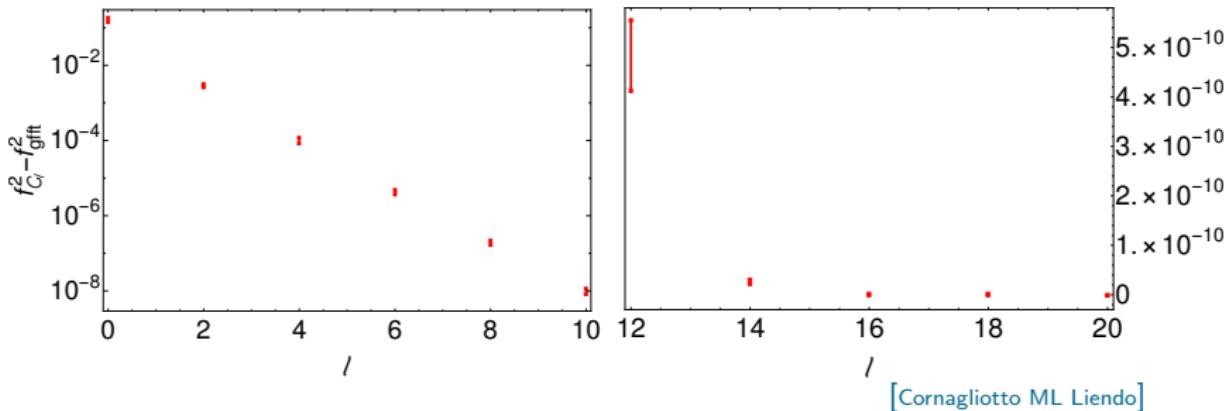
Lorentzian inversion formula

$$\phi\phi \sim f_{\phi^2}^2 \underbrace{\phi^2}_{\Delta=2\Delta_\phi} + f_{C_\ell}^2 \underbrace{C_{\ell>0}}_{\Delta=2\Delta_\phi+\ell} + \dots$$



Lorentzian inversion formula

$$\phi\phi \sim f_{\phi^2}^2 \underbrace{\phi^2}_{\Delta=2\Delta_\phi} + f_{C_\ell}^2 \underbrace{C_{\ell>0}}_{\Delta=2\Delta_\phi+\ell} + \dots$$

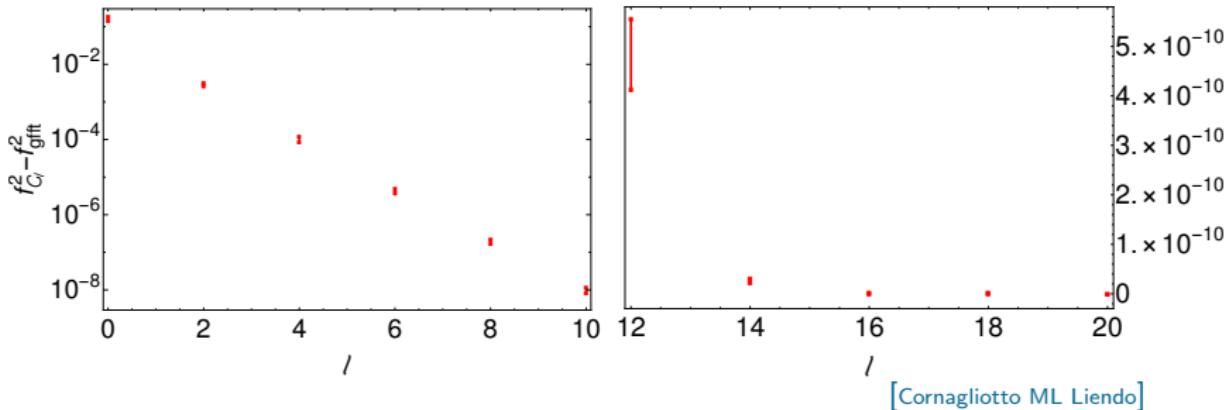


Inverting the $\phi\phi$ OPE

→ Same as bosonic inversion, valid for $\ell > 1$

Lorentzian inversion formula

$$\phi\phi \sim f_{\phi^2}^2 \underbrace{\phi^2}_{\Delta=2\Delta_\phi} + f_{C_\ell}^2 \underbrace{C_{\ell>0}}_{\Delta=2\Delta_\phi+\ell} + \dots$$

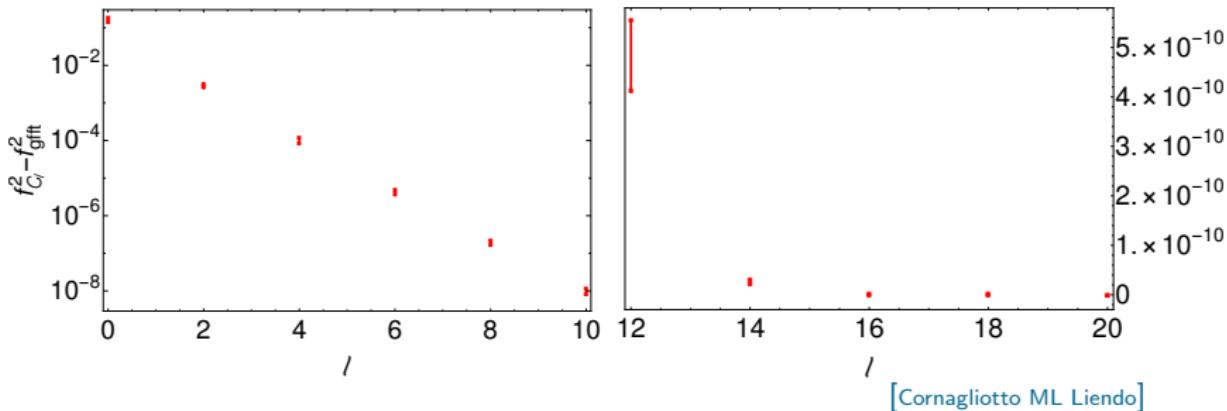


Inverting the $\phi\phi$ OPE

- Same as bosonic inversion, valid for $\ell > 1$
- Feed in low twist in t/u -channel: $\bar{\phi}\phi$ OPE

Lorentzian inversion formula

$$\phi\phi \sim f_{\phi^2}^2 \underbrace{\phi^2}_{\Delta=2\Delta_\phi} + f_{C_\ell}^2 \underbrace{C_{\ell>0}}_{\Delta=2\Delta_\phi+\ell} + \dots$$

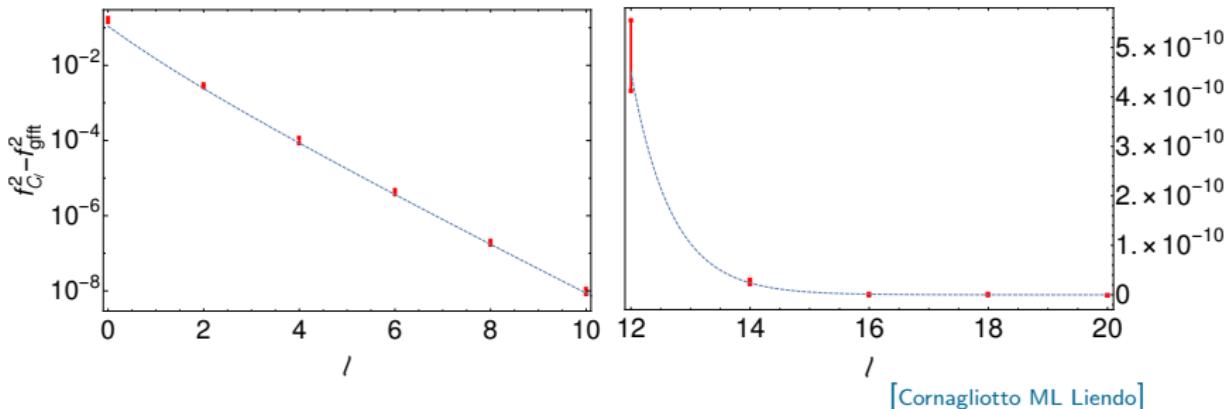


Inverting the $\phi\phi$ OPE

- Same as bosonic inversion, valid for $\ell > 1$
- Feed in low twist in t/u -channel: $\bar{\phi}\phi$ OPE
 - ↪ Only input: $\bar{\phi}\phi \sim 1 + \text{Stress tensor multiplet}$

Lorentzian inversion formula

$$\phi\phi \sim f_{\phi^2}^2 \underbrace{\phi^2}_{\Delta=2\Delta_\phi} + f_{C_\ell}^2 \underbrace{C_{\ell>0}}_{\Delta=2\Delta_\phi+\ell} + \dots$$



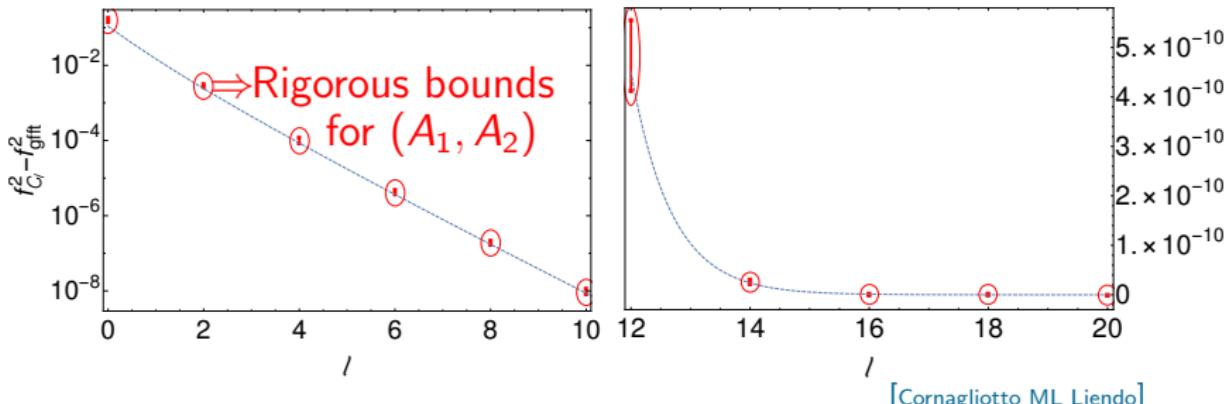
[Cornagliotto ML Liendo]

Inverting the $\phi\phi$ OPE

- Same as bosonic inversion, valid for $\ell > 1$
- Feed in low twist in t/u -channel: $\bar{\phi}\phi$ OPE
 - ↪ Only input: $\bar{\phi}\phi \sim 1 + \text{Stress tensor multiplet}$
- Get s -channel ($\phi\phi$) large spin

Lorentzian inversion formula

$$\phi\phi \sim f_{\phi^2}^2 \underbrace{\phi^2}_{\Delta=2\Delta_\phi} + f_{C_\ell}^2 \underbrace{C_{\ell>0}}_{\Delta=2\Delta_\phi+\ell} + \dots$$

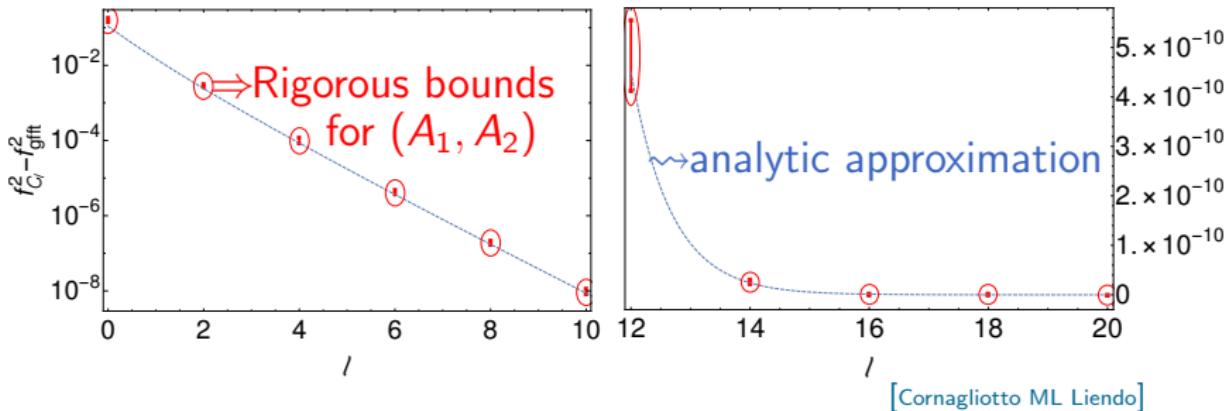


Inverting the $\phi\phi$ OPE

- Same as bosonic inversion, valid for $\ell > 1$
- Feed in low twist in t/u -channel: $\bar{\phi}\phi$ OPE
 - ↪ Only input: $\bar{\phi}\phi \sim 1 + \text{Stress tensor multiplet}$
- Get s -channel ($\phi\phi$) large spin

Lorentzian inversion formula

$$\phi\phi \sim f_{\phi^2}^2 \underbrace{\phi^2}_{\Delta=2\Delta_\phi} + f_{C_\ell}^2 \underbrace{C_{\ell>0}}_{\Delta=2\Delta_\phi+\ell} + \dots$$



Inverting the $\phi\phi$ OPE

- Same as bosonic inversion, valid for $\ell > 1$
- Feed in low twist in t/u -channel: $\bar{\phi}\phi$ OPE
 - ↪ Only input: $\bar{\phi}\phi \sim 1 + \text{Stress tensor multiplet}$
- Get s -channel ($\phi\phi$) large spin

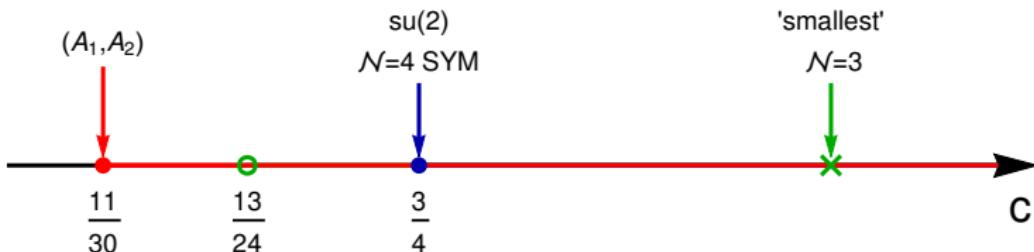
Outline

- ① The Superconformal Bootstrap Program
- ② (A_1, A_2) Argyres-Douglas Theory
- ③ Landscape of $4d \mathcal{N} = 2$ SCFTs
- ④ Summary & Outlook

Landscape of $4d \mathcal{N} \geq 2$ SCFTs

Projection of space of SCFTs to an axis

- $4d \mathcal{N} = 4$ SCFTs $c = a \geq \frac{3}{4}$ [Beem Rastelli van Rees]
- $4d \mathcal{N} \geq 3$ SCFTs $c = a > \frac{13}{24}$ [Cornagliotto ML Schomerus]
- $4d \mathcal{N} \geq 2$ SCFTs $c \geq \frac{11}{30}$ [Liendo Ramirez Seo]



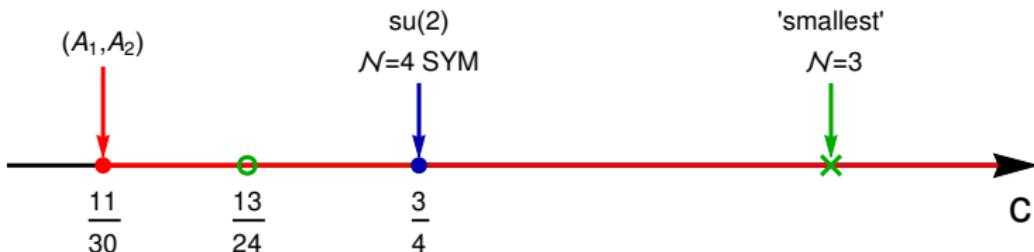
Landscape of $4d \mathcal{N} \geq 2$ SCFTs

Projection of space of SCFTs to an axis

→ $4d \mathcal{N} = 4$ SCFTs $c = a \geq \frac{3}{4}$ [Beem Rastelli van Rees]

→ $4d \mathcal{N} \geq 3$ SCFTs $c = a > \frac{13}{24}$ [Cornagliotto ML Schomerus]

→ $4d \mathcal{N} \geq 2$ SCFTs $c \geq \frac{11}{30}$ [Liendo Ramirez Seo]



Finer view of the space of theories:

⇒ Organize theories by flavor symmetry

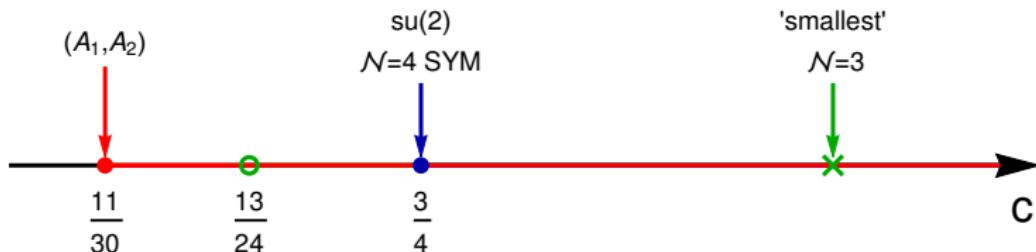
Landscape of $4d \mathcal{N} \geq 2$ SCFTs

Projection of space of SCFTs to an axis

→ $4d \mathcal{N} = 4$ SCFTs $c = a \geq \frac{3}{4}$ [Beem Rastelli van Rees]

→ $4d \mathcal{N} \geq 3$ SCFTs $c = a > \frac{13}{24}$ [Cornagliotto ML Schomerus]

→ $4d \mathcal{N} \geq 2$ SCFTs $c \geq \frac{11}{30}$ [Liendo Ramirez Seo]



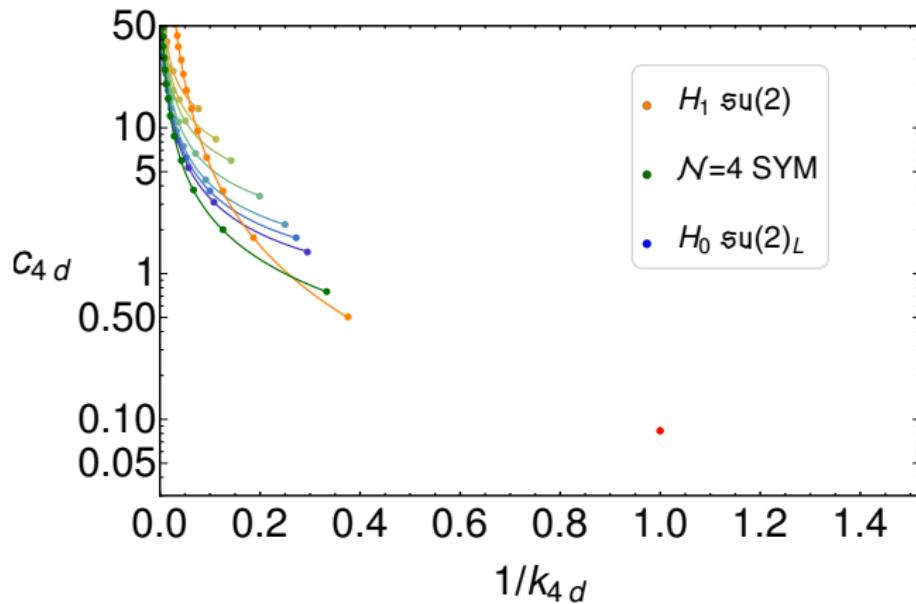
Finer view of the space of theories:

⇒ Organize theories by flavor symmetry

$$\langle TT \rangle \propto c, \quad \langle JJ \rangle \propto k$$

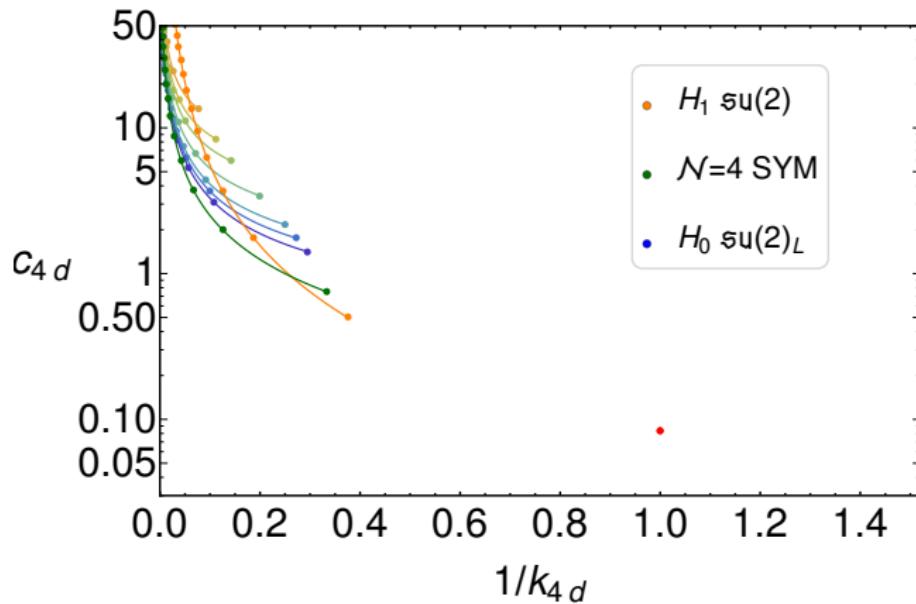
$4d \mathcal{N} = 2$ SCFT with $su(2)$ flavor symmetry

- ▶ 4d Flavor current supermultiplet



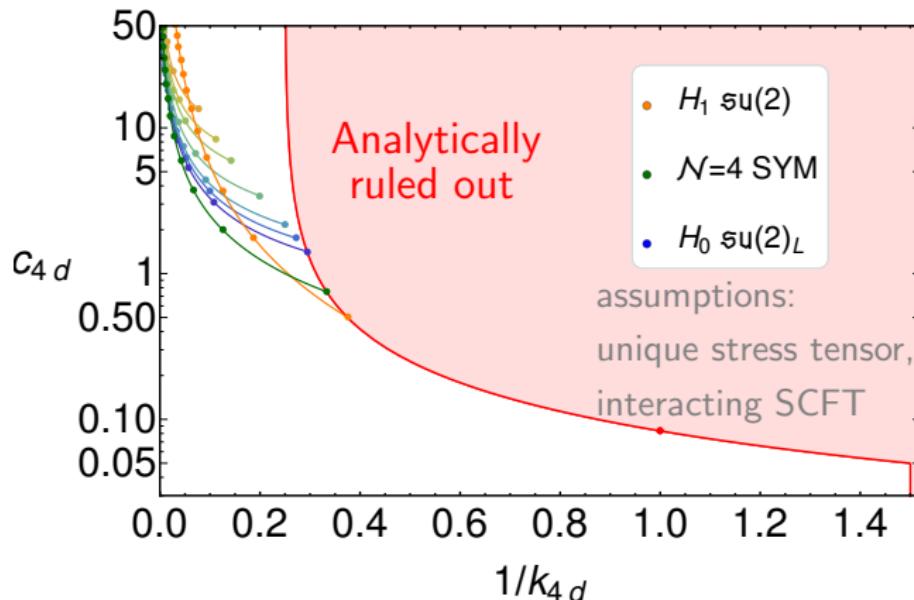
$4d \mathcal{N} = 2$ SCFT with $su(2)$ flavor symmetry

- ▶ $4d$ Flavor current supermultiplet $\mapsto \langle JJJJ \rangle_{2d}$



$4d \mathcal{N} = 2$ SCFT with $su(2)$ flavor symmetry

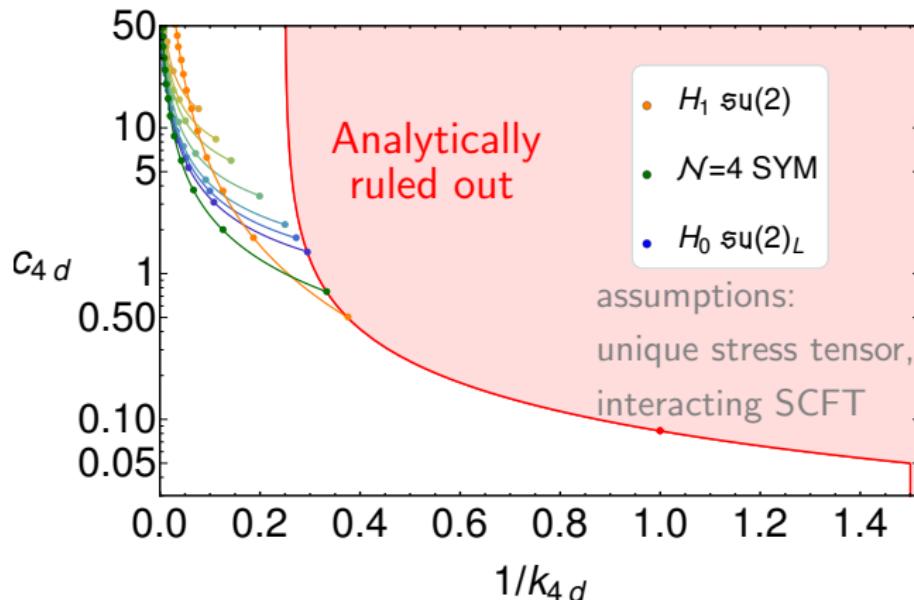
- ▶ $4d$ Flavor current supermultiplet $\mapsto \langle JJJJ \rangle_{2d} \rightsquigarrow f_{4d}^2 \geq 0$



[Beem ML Liendo Peelaers Rastelli van Rees]

$4d \mathcal{N} = 2$ SCFT with $su(2)$ flavor symmetry

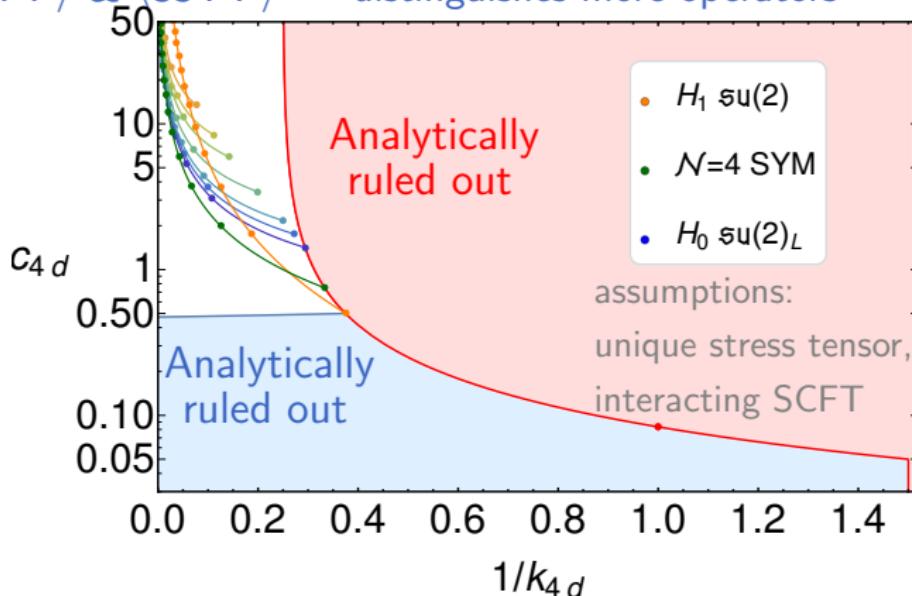
- ▶ $4d$ Flavor current supermultiplet $\mapsto \langle JJJJ \rangle_{2d} \rightsquigarrow \sum f_{4d}^2 \geq 0$



[Beem ML Liendo Peelaers Rastelli van Rees]

$4d \mathcal{N} = 2$ SCFT with $su(2)$ flavor symmetry

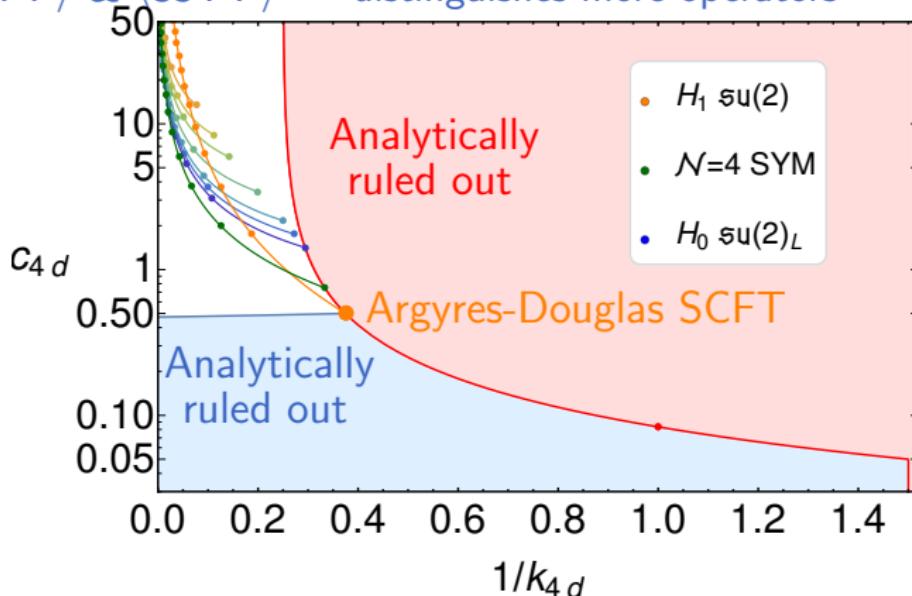
- ▶ 4d Flavor current supermultiplet $\mapsto \langle JJJJ \rangle_{2d} \rightsquigarrow \sum f_{4d}^2 \geq 0$
- ▶ $\langle TTTT \rangle$ & $\langle JJTT \rangle \rightsquigarrow$ distinguishes more operators



[Beem ML Liendo Peelaers Rastelli van Rees, ML Liendo]

$4d \mathcal{N} = 2$ SCFT with $su(2)$ flavor symmetry

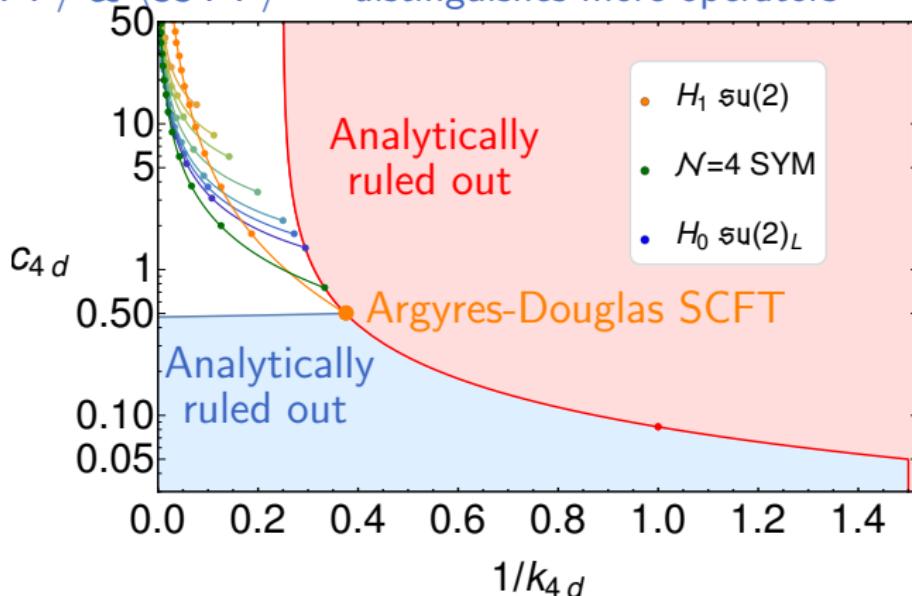
- ▶ $4d$ Flavor current supermultiplet $\mapsto \langle JJJJ \rangle_{2d} \rightsquigarrow \sum f_{4d}^2 \geq 0$
- ▶ $\langle TTTT \rangle$ & $\langle JJTT \rangle \rightsquigarrow$ distinguishes more operators



[Beem ML Liendo Peelaers Rastelli van Rees, ML Liendo]

$4d \mathcal{N} = 2$ SCFT with $su(2)$ flavor symmetry

- ▶ $4d$ Flavor current supermultiplet $\mapsto \langle JJJJ \rangle_{2d} \rightsquigarrow \sum f_{4d}^2 \geq 0$
- ▶ $\langle TTTT \rangle$ & $\langle JJTT \rangle \rightsquigarrow$ distinguishes more operators



Only for $su(2)$, $su(3)$, $so(8)$, g_2 , f_4 , e_6 , e_7 , e_8

[Beem ML Liendo Peelaers Rastelli van Rees, ML Liendo]

Outline

- ① The Superconformal Bootstrap Program
- ② (A_1, A_2) Argyres-Douglas Theory
- ③ Landscape of $4d \mathcal{N} = 2$ SCFTs
- ④ Summary & Outlook

Summary & Outlook

Constrained the “simplest” Argyres-Douglas theory

Summary & Outlook

Constrained the “simplest” Argyres-Douglas theory

Zoom in to other strongly coupled $\mathcal{N} = 2$ SCFTs?

(at corners of $su(2)$, $su(3)$, e_6 , e_7 , e_8 exclusion curves)

Summary & Outlook

Constrained the “simplest” Argyres-Douglas theory

Zoom in to other strongly coupled $\mathcal{N} = 2$ SCFTs?

(at corners of $su(2)$, $su(3)$, e_6 , e_7 , e_8 exclusion curves)

→ Mixed system: stress tensor & flavor current multiplets

Summary & Outlook

Constrained the “simplest” Argyres-Douglas theory

Zoom in to other strongly coupled $\mathcal{N} = 2$ SCFTs?

(at corners of $su(2)$, $su(3)$, e_6 , e_7 , e_8 exclusion curves)

- Mixed system: stress tensor & flavor current multiplets
- Stronger numerical constraints on the space of theories?

Summary & Outlook

Constrained the “simplest” Argyres-Douglas theory

Zoom in to other strongly coupled $\mathcal{N} = 2$ SCFTs?

(at corners of $su(2)$, $su(3)$, e_6 , e_7 , e_8 exclusion curves)

- Mixed system: stress tensor & flavor current multiplets
- Stronger numerical constraints on the space of theories?

Superblocks for Super-stress tensor multiplets

Summary & Outlook

Constrained the “simplest” Argyres-Douglas theory

Zoom in to other strongly coupled $\mathcal{N} = 2$ SCFTs?

(at corners of $su(2)$, $su(3)$, e_6 , e_7 , e_8 exclusion curves)

- Mixed system: stress tensor & flavor current multiplets
- Stronger numerical constraints on the space of theories?

Superblocks for Super-stress tensor multiplets

- Bounds on (c, k) did not come from superprimary of stress tensor

Summary & Outlook

Constrained the “simplest” Argyres-Douglas theory

Zoom in to other strongly coupled $\mathcal{N} = 2$ SCFTs?

(at corners of $su(2)$, $su(3)$, e_6 , e_7 , e_8 exclusion curves)

- Mixed system: stress tensor & flavor current multiplets
- Stronger numerical constraints on the space of theories?

Superblocks for Super-stress tensor multiplets

- Bounds on (c, k) did not come from superprimary of stress tensor – compute whole superblock?

Summary & Outlook

Constrained the “simplest” Argyres-Douglas theory

Zoom in to other strongly coupled $\mathcal{N} = 2$ SCFTs?

(at corners of $su(2)$, $su(3)$, e_6 , e_7 , e_8 exclusion curves)

- Mixed system: stress tensor & flavor current multiplets
- Stronger numerical constraints on the space of theories?

Superblocks for Super-stress tensor multiplets

- Bounds on (c, k) did not come from superprimary of stress tensor – compute whole superblock?
- Two-dimensional long blocks [Cornagliotto ML Schomerus]
needed for $c > \frac{13}{24}$ for $\mathcal{N} = 3$ SCFTs

Summary & Outlook

Constrained the “simplest” Argyres-Douglas theory

Zoom in to other strongly coupled $\mathcal{N} = 2$ SCFTs?

(at corners of $su(2)$, $su(3)$, e_6 , e_7 , e_8 exclusion curves)

- Mixed system: stress tensor & flavor current multiplets
- Stronger numerical constraints on the space of theories?

Superblocks for Super-stress tensor multiplets

- Bounds on (c, k) did not come from superprimary of stress tensor – compute whole superblock?
- Two-dimensional long blocks [Cornagliotto ML Schomerus]
needed for $c > \frac{13}{24}$ for $\mathcal{N} = 3$ SCFTs
- Weight-shifting operators? [Karateev Kravchuk Simmons-Duffin]

Summary & Outlook

Constrained the “simplest” Argyres-Douglas theory

Zoom in to other strongly coupled $\mathcal{N} = 2$ SCFTs?

(at corners of $su(2)$, $su(3)$, e_6 , e_7 , e_8 exclusion curves)

- Mixed system: stress tensor & flavor current multiplets
- Stronger numerical constraints on the space of theories?

Superblocks for Super-stress tensor multiplets

- Bounds on (c, k) did not come from superprimary of stress tensor – compute whole superblock?
- Two-dimensional long blocks [Cornagliotto ML Schomerus]
needed for $c > \frac{13}{24}$ for $\mathcal{N} = 3$ SCFTs
- Weight-shifting operators? [Karateev Kravchuk Simmons-Duffin]

What is the “smallest” $\mathcal{N} = 3$ SCFT?

Thank you!

Backup slides

Outline

- ⑤ Lorentzian inversion formula for (A_1, A_2)
- ⑥ Constraining the space of $4d \mathcal{N} = 2$ SCFTs

Lorentzian inversion formula: Superconformal case

Invert $\phi\phi$ OPE

- Same as bosonic inversion, valid for $\ell > 1$
- Feed in $\bar{\phi}\phi \sim 1 + \text{Stress tensor multiplet} + \dots$

Lorentzian inversion formula: Superconformal case

Invert $\phi\phi$ OPE

- Same as bosonic inversion, valid for $\ell > 1$
- Feed in $\bar{\phi}\phi \sim 1 + \text{Stress tensor multiplet} + \dots$

Invert $\bar{\phi}\phi$ OPE

- Supersymmetric inversion: valid for $\ell \geq 0$

Lorentzian inversion formula: Superconformal case

Invert $\phi\phi$ OPE

- Same as bosonic inversion, valid for $\ell > 1$
- Feed in $\bar{\phi}\phi \sim 1 + \text{Stress tensor multiplet} + \dots$

Invert $\bar{\phi}\phi$ OPE

- Supersymmetric inversion: valid for $\ell \geq 0$
- Feed in low twist in t -channel ($\bar{\phi}\phi$)

Lorentzian inversion formula: Superconformal case

Invert $\phi\phi$ OPE

- Same as bosonic inversion, valid for $\ell > 1$
- Feed in $\bar{\phi}\phi \sim 1 + \text{Stress tensor multiplet} + \dots$

Invert $\bar{\phi}\phi$ OPE

- Supersymmetric inversion: valid for $\ell \geq 0$
- Feed in low twist in t -channel ($\bar{\phi}\phi$)
 - $\bar{\phi}\phi \sim 1 + \text{Stress tensor multiplet} + \dots$

Lorentzian inversion formula: Superconformal case

Invert $\phi\phi$ OPE

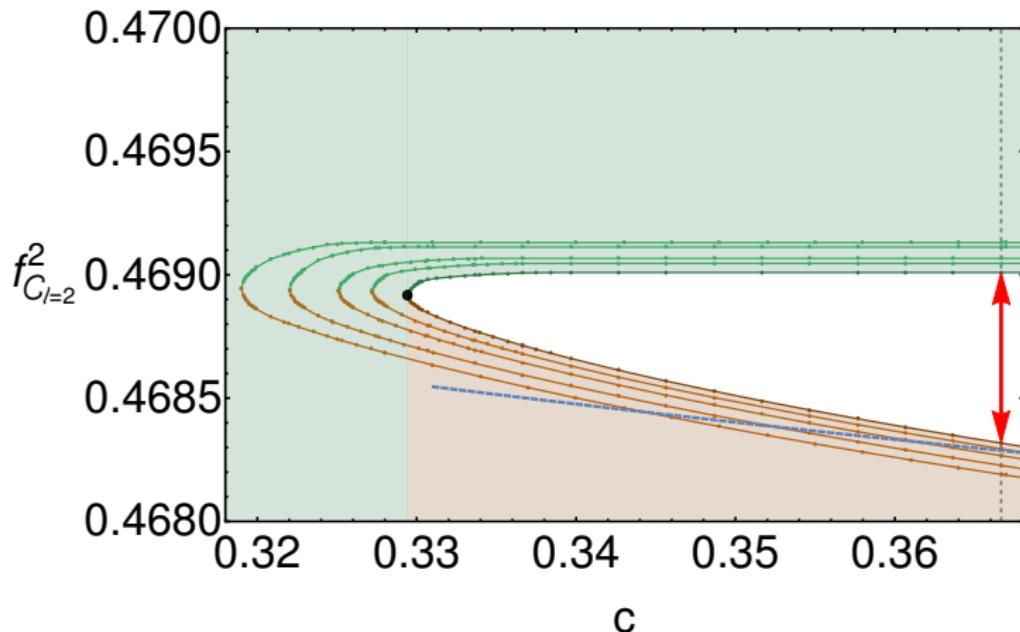
- Same as bosonic inversion, valid for $\ell > 1$
- Feed in $\bar{\phi}\phi \sim 1 + \text{Stress tensor multiplet} + \dots$

Invert $\bar{\phi}\phi$ OPE

- Supersymmetric inversion: valid for $\ell \geq 0$
- Feed in low twist in t -channel ($\bar{\phi}\phi$)
 - $\bar{\phi}\phi \sim 1 + \text{Stress tensor multiplet} + \dots$
- and in u -channel ($\phi\phi$)
 - $\phi\phi \sim \phi^2 + \dots$

Bounding OPE coefficients

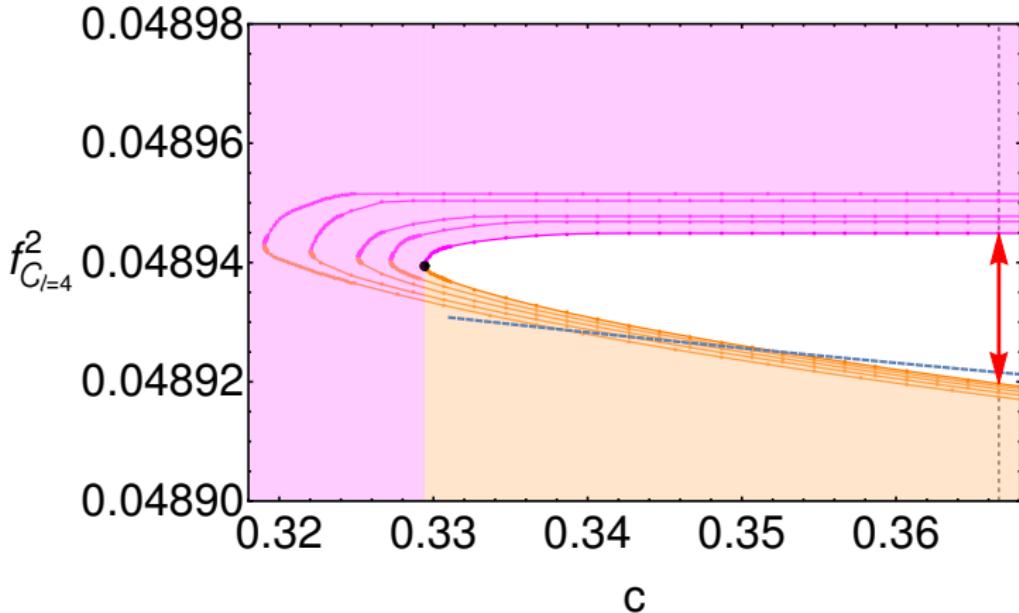
$$\phi\phi \sim f_{\phi^2}^2 \underbrace{\phi^2}_{\Delta=2\Delta_\phi} + f_{C_\ell}^2 \underbrace{C_{\ell>0}}_{\Delta=2\Delta_\phi+\ell} + \dots$$



[Cornagliotto ML Liendo]

Bounding OPE coefficients

$$\phi\phi \sim f_{\phi^2}^2 \underbrace{\phi^2}_{\Delta=2\Delta_\phi} + f_{C_\ell}^2 \underbrace{C_{\ell>0}}_{\Delta=2\Delta_\phi+\ell} + \dots$$



[Cornagliotto ML Liendo]

A Lorentzian inversion formula

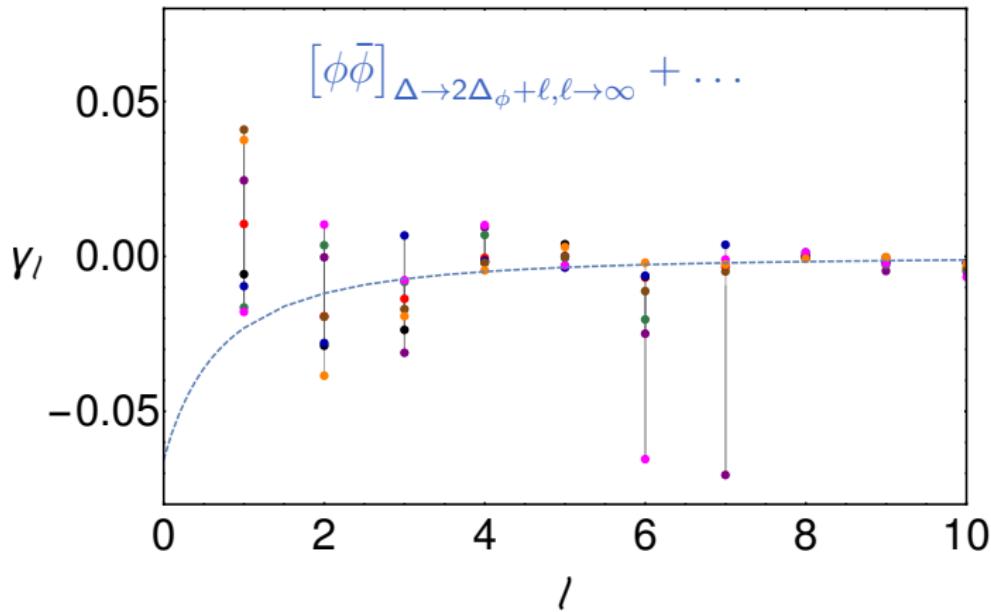
Inverting the $\phi\bar{\phi}$ OPE

- Supersymmetric inversion: valid for $\ell \geq 0$
- Only input: $\bar{\phi}\phi \sim 1 + \text{Stress tensor multiplet}$

A Lorentzian inversion formula

Inverting the $\phi\bar{\phi}$ OPE

- Supersymmetric inversion: valid for $\ell \geq 0$
- Only input: $\bar{\phi}\phi \sim 1 + \text{Stress tensor multiplet}$



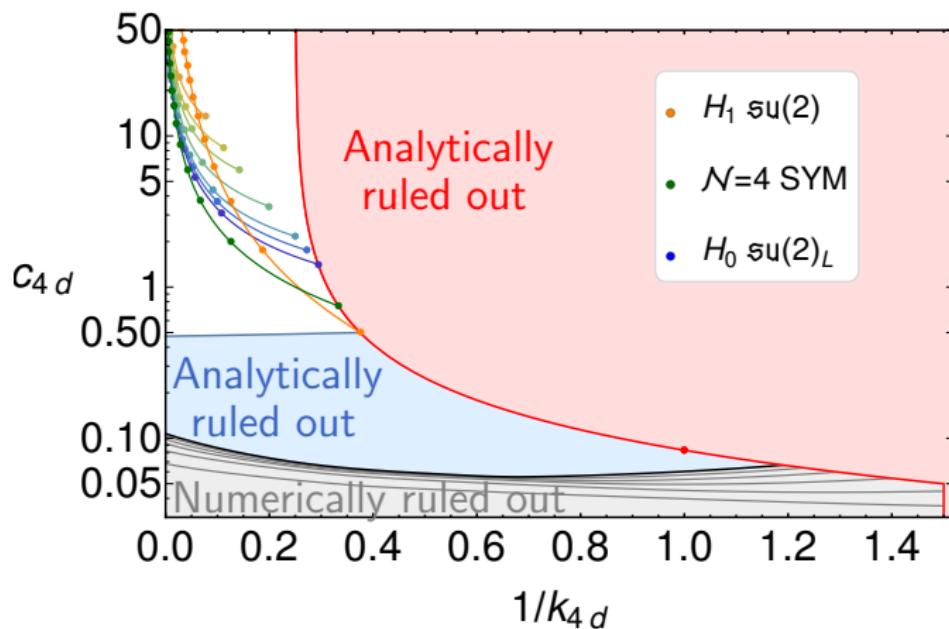
Outline

- ⑤ Lorentzian inversion formula for (A_1, A_2)

- ⑥ Constraining the space of $4d \mathcal{N} = 2$ SCFTs

Constraining the space of $4d \mathcal{N} = 2$ SCFTs

$su(2)$ flavor symmetry

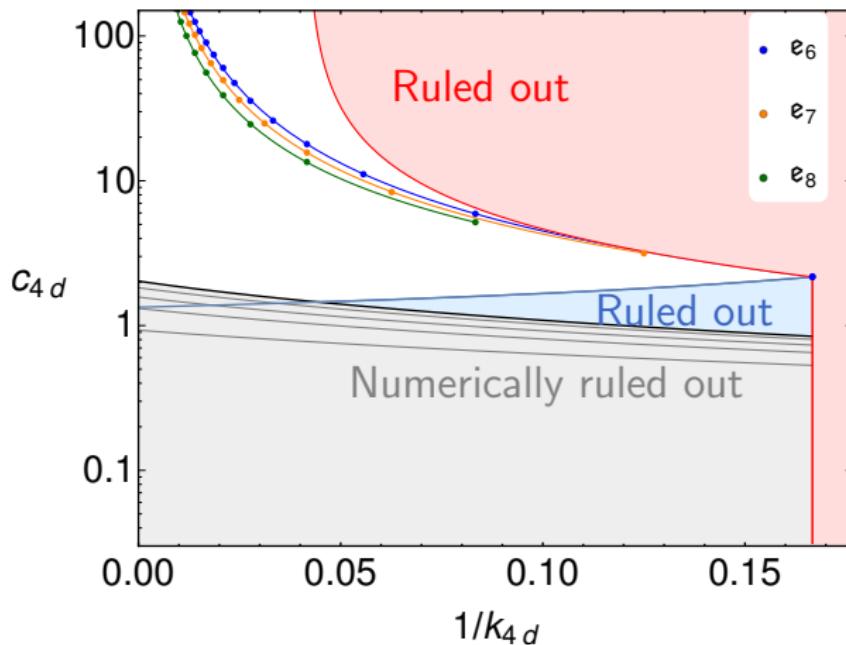


[Beem, ML, Liendo, Peelaers, Rastelli, van Rees; ML, Liendo]

[Beem, ML, Liendo, Rastelli, van Rees]

Constraining the space of $4d \mathcal{N} = 2$ SCFTs

e_6 flavor symmetry

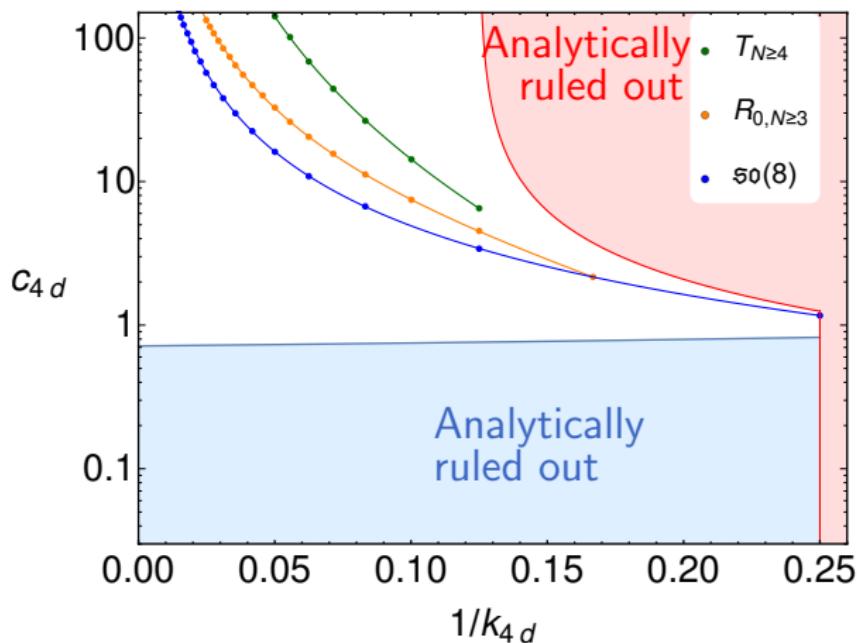


[Beem, ML, Liendo, Peelaers, Rastelli, van Rees; ML, Liendo]

[Beem, ML, Liendo, Rastelli, van Rees]

Constraining the space of $4d \mathcal{N} = 2$ SCFTs

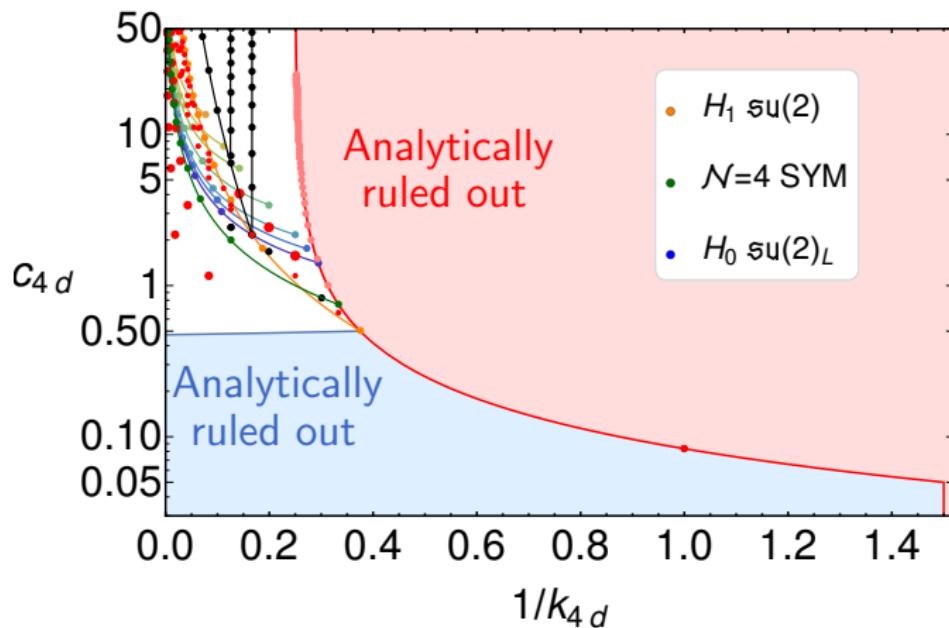
$su(4)$ flavor symmetry



[Beem, ML, Liendo, Peelaers, Rastelli, van Rees; ML, Liendo]

Constraining the space of $4d \mathcal{N} = 2$ SCFTs

$su(2)$ flavor symmetry



[Beem, ML, Liendo, Peelaers, Rastelli, van Rees; ML, Liendo]